

Small angle neutron scattering

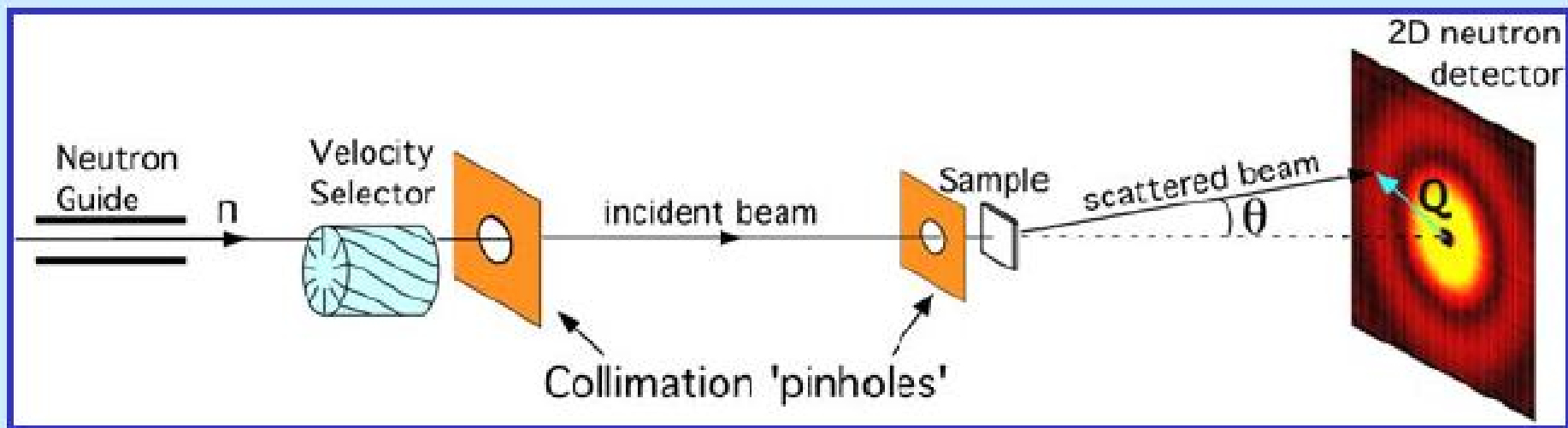
Focus on magnetic nanoparticles and biological applications

A.V.Belushkin

**Frank Laboratory of Neutron Physics, JINR, Dubna,
Russia**



Small Angle Neutron Scattering (SANS)



Small-Angle Neutron Scattering (SANS) probes structure on a scale d , where

$$d \approx \frac{\lambda}{\theta}$$

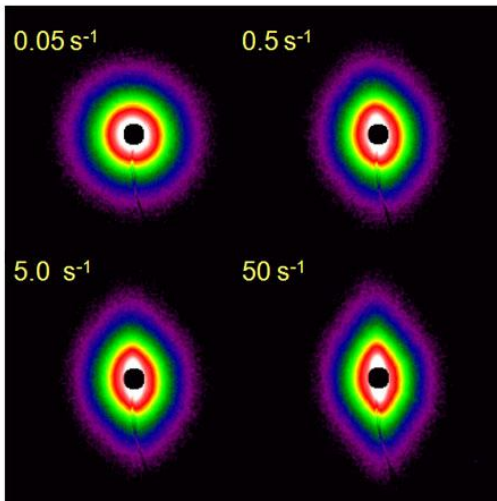
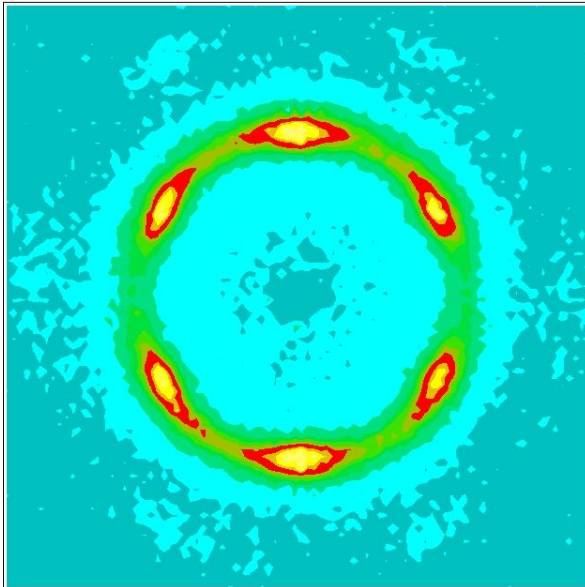
(wavelength)
(scattering angle)

$0.5 \text{ nm} < \lambda < 2 \text{ nm}$ (cold neutrons)

$0.1^\circ < \theta < 10^\circ$ (small angles)

$1 \text{ nm} < d < 300 \text{ nm}$

Examples of registered patterns on a detector



From the analysis of the signal on a detector, using advanced mathematical models one can extract information about characteristic sizes, shape and relative orientation as well as some other parameters of the nanoparticles in a sample under investigation

Analyze Solution Scattering Automatically

Style: Frame Facet Gouraud Function: Show Hide Delete Config... File... Command ▾ Load... Active ▾ No 9. shout Quit

Active structure attributes
Type: Single Edges: Invisible
Colour: aquamarine HLHSR: Enable
Visible: All Transp. %: 25
Annotation: 7

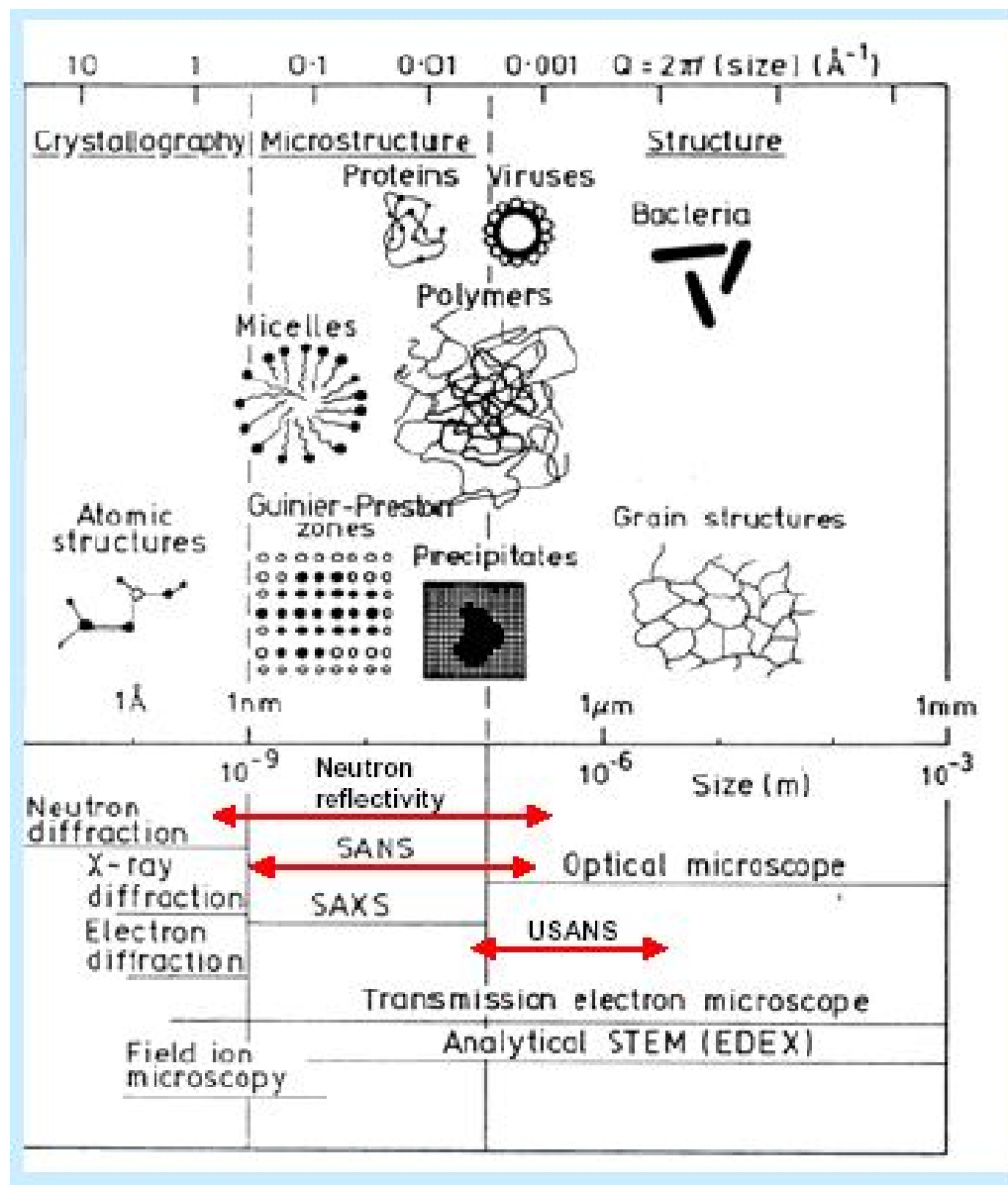
TRANSFORMATION: Man Def Reset
Rotate X: -53 -180 180
Rotate Y: 54 -180 180
Rotate Z: 0 -180 180
Shift X: 0 -210 210
Shift Y: 0 -210 210
Shift Z: 0 -315 315
Scale (%): 100 10 200
Mode: Discrete Kbd Delta (*PI/100): 10

Shape determination
Input intensity file [dat]: sacctet
Angular units (1-4): 2
Fitting range (fractional): 1.0
Output file [flm]: shout
Init. approximation [flm]:
Maximum order of harmonics: 4
Apply the filter: N
Go Stop Resume Cancel

Reper: Visible MESSAGE: shout is loaded
Length: 20

MKPlot
s, I(s)exp, I(s)calc for shout.flm12-JUN-1996
Wed Jun 12 18:57:00 1996
Y
10.5
10.0
9.5
9.0
8.5
8.0
7.5
7.0
X
0.05 0.10 0.15 0.20 0.25 0.30

SANS applicability



Основные теоретические представления

При рассмотрении малоуглового рассеяния нейтронов вводится понятие пространственного распределения плотности когерентной амплитуды рассеяния

$$\rho(\vec{r}) = \bar{b} \times N_a(\vec{r})$$

где b - средняя когерентная амплитуда рассеяния нейтронов для частицы (или соответственно растворителя), N_a - распределение атомной плотности в частице (или растворителе), имеющее размерность [число атомов / единица объема]. Потенциал взаимодействия нейтрона с таким объектом может быть записан в виде

$$H_c = \frac{2 \pi \hbar^2}{m} \rho(\vec{r})$$

При рассеянии плоской нейтронной волны на таком потенциале амплитуда упругого и когерентного рассеяния будет пропорциональна матричному элементу

$$\langle \vec{k}_1 | H_c | \vec{k}_0 \rangle$$



$$\langle \vec{k}_1 | H_c | \vec{k}_0 \rangle = A(\vec{Q}) = \int_V \rho(\vec{r}) \exp(i\vec{Q}\vec{r}) d\vec{r}$$

где интегрирование идет по всему объему исследуемого образца. Интенсивность рассеянной волны пропорциональна дифференциальному сечению рассеяния, т.е. квадрату матричного элемента.

$$I(\vec{Q}) \sim \frac{d\sigma}{d\Omega} = |A(\vec{Q})|^2 = \left| \int_V \rho(\vec{r}) \exp(i\vec{Q}\vec{r}) d\vec{r} \right|^2$$

Задача исследования методом малоуглового рассеяния состоит в восстановлении функции $\rho(\underline{r})$ в исследуемом объекте по измеренной зависимости $I(\underline{Q})$. Степень разрешимости этой задачи зависит от многих факторов. Рассмотрим некоторые наиболее часто используемые приближения.



Изотропное распределение частиц в растворе

Выражение для сечения рассеяния следует усреднить по всем возможным пространственным ориентациям частиц относительно вектора рассеяния:

$$\left\langle \frac{d\sigma}{d\Omega} \right\rangle = \int_V \int_{V'} \rho(\vec{r}) \rho(\vec{r}') \frac{\sin Q|\vec{r} - \vec{r}'|}{Q|\vec{r} - \vec{r}'|} d\vec{r} d\vec{r}'$$

Формула носит название формулы Дебая. Она получается на основе следующих соотношений

$$\begin{aligned} \langle \sin[\vec{Q}(\vec{r} - \vec{r}')] \rangle &= \int_0^\pi \sin(Q|\vec{r} - \vec{r}'| \cos \varphi) \frac{\sin \varphi}{2} d\varphi = 0 \\ \langle \cos[\vec{Q}(\vec{r} - \vec{r}')] \rangle &= \int_0^\pi \cos(Q|\vec{r} - \vec{r}'| \cos \varphi) \frac{\sin \varphi}{2} d\varphi = \frac{\sin Q|\vec{r} - \vec{r}'|}{Q|\vec{r} - \vec{r}'|} \end{aligned}$$

Здесь φ - угол между векторами \underline{Q} и $\underline{r-r'}$, а $\sin(\varphi)d\varphi/2$ - вероятность того, что угол φ заключен в интервале от φ до $\varphi + d\varphi$.



Во всех вышеприведенных рассуждениях мы считали, что частицы находятся в вакууме, т.е. не учитывали свойств растворителя. Все вышеприведенные формулы при учете растворителя остаются справедливыми, если заменить ρ на $\rho - \rho_s$, где ρ_s - плотность когерентного рассеяния растворителя. При этом вводится важное понятие **контраста**.

Контраст – разницы между средней плотностью рассеяния частицы и плотностью рассеяния растворителя:

$$\Delta\rho = \langle \rho(\vec{r}) - \rho_s \rangle = \frac{1}{V} \int [\rho(\vec{r}) - \rho_s] d\vec{r} = \bar{\rho} - \rho_s$$

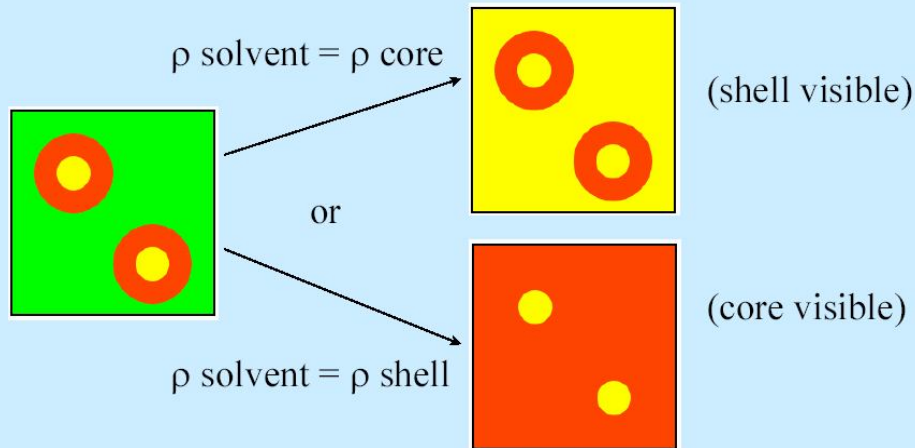
При больших значениях $\Delta\rho$ (сильный контраст) сигнал от частиц будет велик по сравнению с сигналом от растворителя, т.е. частицы будут ясно "видны" в малоугловом рассеянии. Наоборот, при слабом контрасте сигнал от частиц будет практически неотделим от общего сигнала. Поэтому контраст является одной из важнейших характеристик в малоугловых экспериментах. **Вариация контраста**, т.е. изменение значения $\Delta\rho$, например, за счет изотопного замещения служит мощным инструментом получения дополнительной структурной информации.



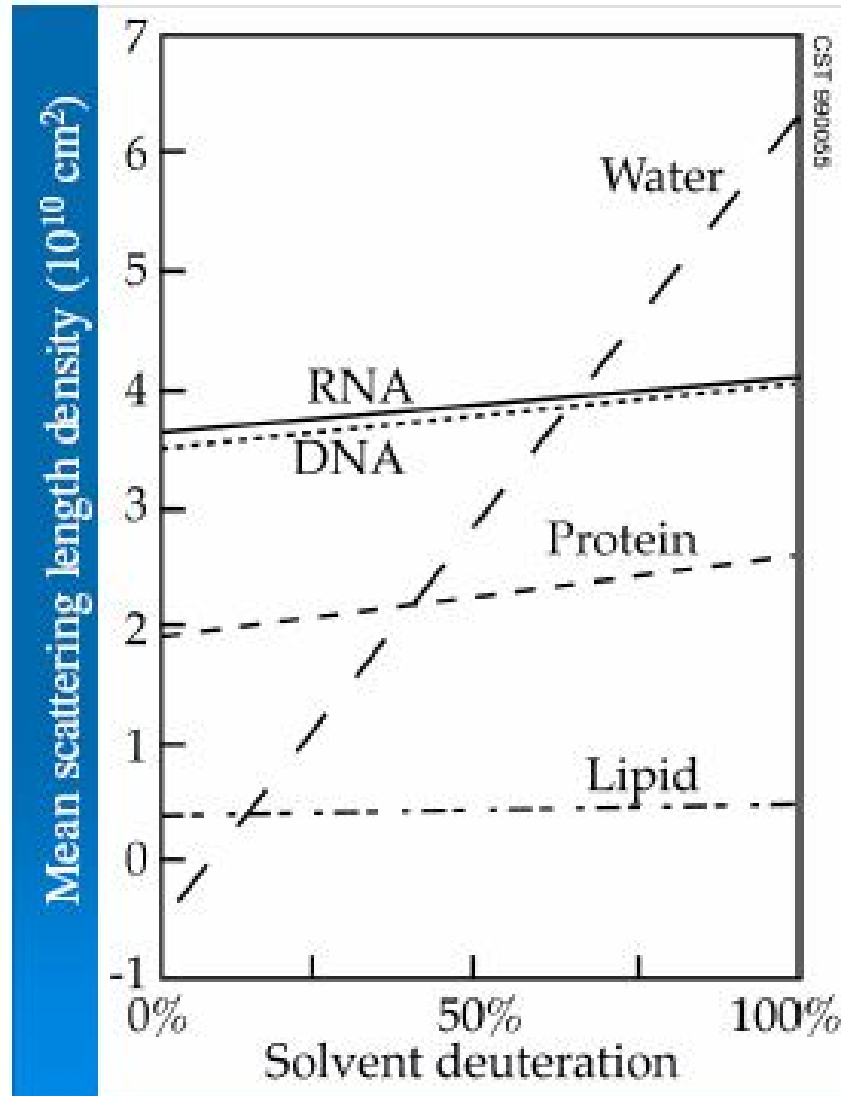
Схематическая иллюстрация вариации контраста

Contrast Matching

reduce the number of phases “visible”



- The two distinct two - phase systems can be easily understood



Guinier approximation

in 1939 A.Guinier simplified Debye formula for $QL \ll 1$. using an expansion $\sin x/x = 1 - x^2/3! + x^4/5! - \dots$ And keeping only the first two terms, Debye formula becomes

$$I(Q) \sim \left\langle \frac{d\sigma}{d\Omega} \right\rangle = \left(\int_V \rho(\vec{r}) d\vec{r} \right)^2 \left(1 - \frac{1}{3} Q^2 \frac{\int_V \rho(\vec{r}) r^2 d\vec{r}}{\int_V \rho(\vec{r}) d\vec{r}} + \dots \right) =$$
$$= \left(\int_V \rho(\vec{r}) d\vec{r} \right)^2 \left((1 - Q^2 R_g^2 / 3 + \dots) \right),$$

where R_g^2 is the radius of gyration with the replacement of the body mass density for scattering density.



Guinier assumed that the expression in last parenthesis can be written in a form

$$F^2(Q) = \exp(-Q^2 R_g^2 / 3)$$

This expression is the basis for extracting the particle radius of gyration from the experimental data. For this purpose the data are plotted as the so called Guinier plot - $\ln I(Q^2)$ and inclination angle gives us R_g :

$$R_g^2 = -\frac{1}{3} d \ln I(Q^2) / dQ^2$$



Forward scattering

For $Q = 0$ the differential cross section becomes

$$\left\langle \frac{d\sigma}{d\Omega} \right\rangle = (\bar{\rho} - \rho_s)^2 V^2$$

By using $I(0)$ one can estimate both particle volume and its molecular mass. For molecular mass estimation one needs the knowledge of average coherent scattering length density and an independently measured a partial specific particle volume

$$\bar{v} = \frac{VN_A}{M}$$

where N_A – is the Avogadro number. Then one has

$$I(0) = N \left\langle \frac{d\sigma}{d\Omega} \right\rangle = N(\bar{\rho} - \rho_s)^2 V^2 = Mc(\bar{\rho} - \rho_s)^2 \bar{v}^2 V_s \cdot N_A^{-1}$$

where N – is the total number of particles in solution, which can be calculated from their weight concentration C (g mole/l) $(M \cdot N)/(V_s/N_A)$, V_s – solvent volume.



Porod integral

One of the important integral characteristic of SANS is the so-called Porod invariant

$$A = \int_0^{\infty} Q^2 I(Q) dQ = 2\pi^2 \int_V [\rho(\vec{r}) - \rho_s]^2 d\vec{r}$$

Porod invariant characterizes the total scattering power of the object under investigation and is proportional to the square of the particle contrast with respect of the solvent. It is used for calculations of different structural parameters. In the case of uniform density particle, Porod invariant directly gives us a particle volume

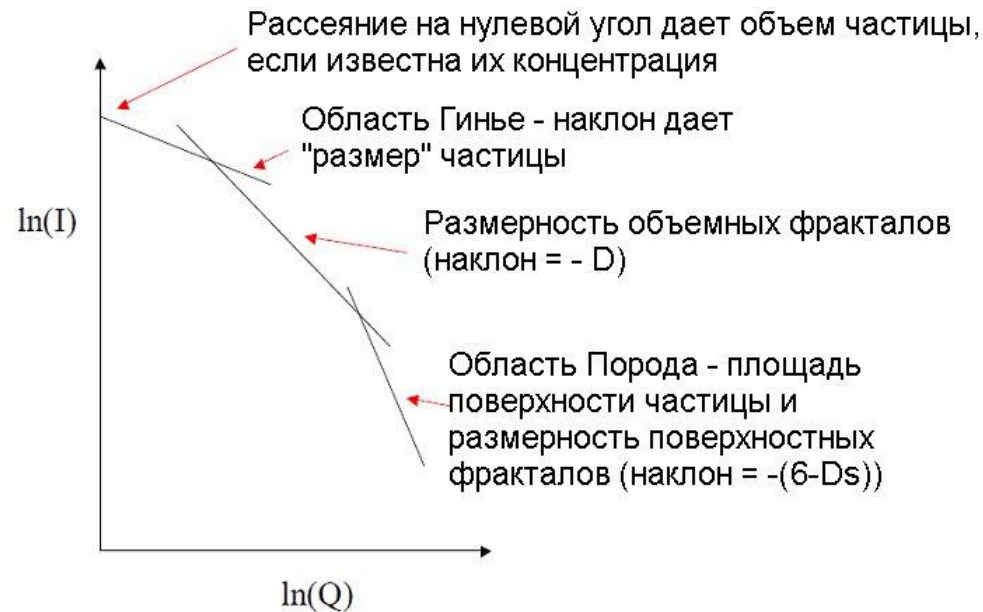
$$A = 2\pi^2 (\bar{\rho} - \rho_s)^2 \cdot V$$



In the limit $Q \rightarrow \infty$ for the uniform density particles one gets

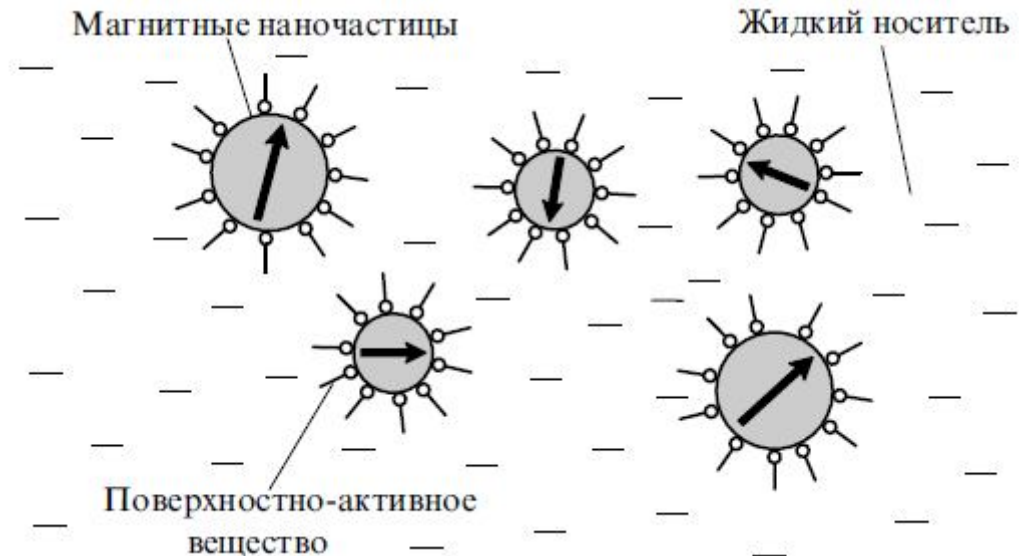
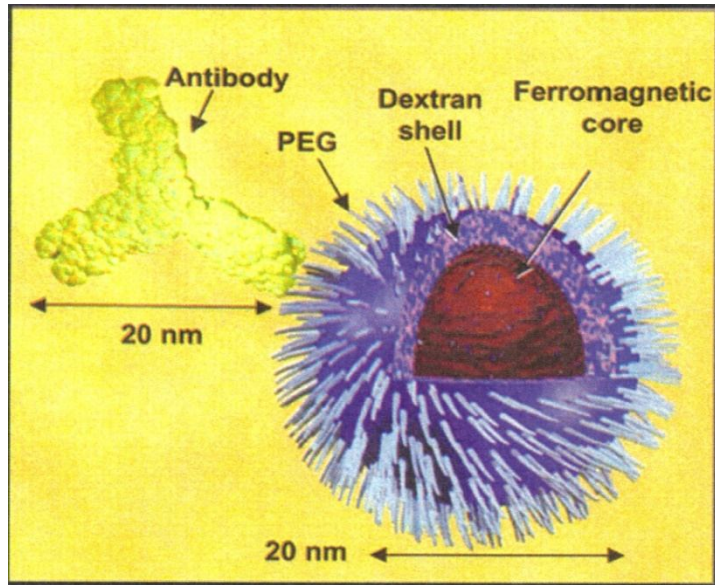
$$I(Q) \underset{Q \rightarrow \infty}{=} \frac{2\pi}{Q^4} (\bar{\rho} - \rho_s)^2 S$$

where S – is the surface area of a particle. Therefore, for the uniform particles it is possible to extract information about their surface area



Information which can be extracted from different parts of the SANS curve

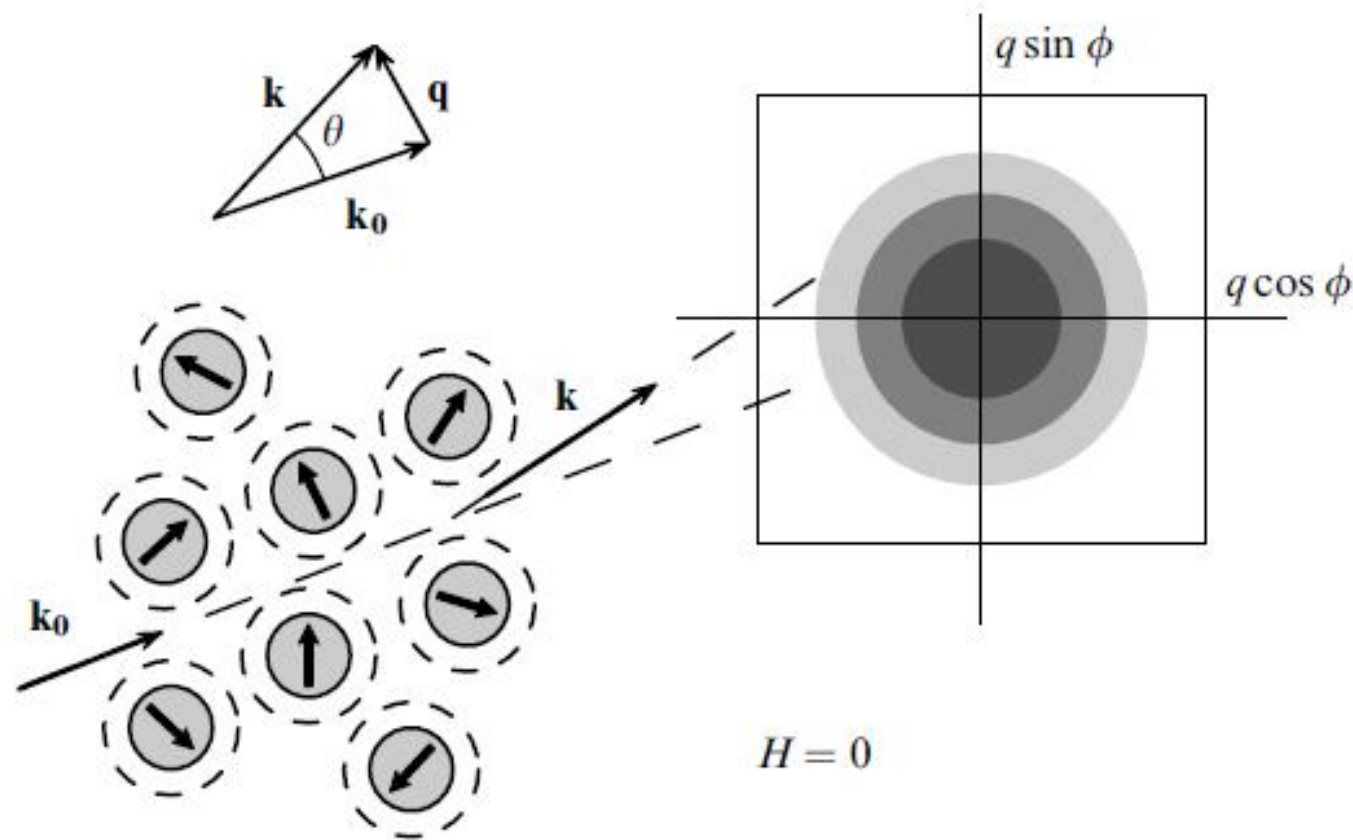
Small Angle Neutron Scattering (SANS) from magnetic nanoparticles in solution



Scattering from a single nanoparticle:

- Nuclear scattering from magnetic core
- Nuclear scattering from hydrogen containing shell
- Magnetic scattering from magnetic core
- Nuclear scattering from the solvent

SANS from magnetic nanoparticles in solution



Principal scheme of SANS on magnetic nanoparticles in solution. In the absence of external magnetic field magnetic moments of nanoparticles are randomly oriented in space.

SANS from magnetic nanoparticles in solution

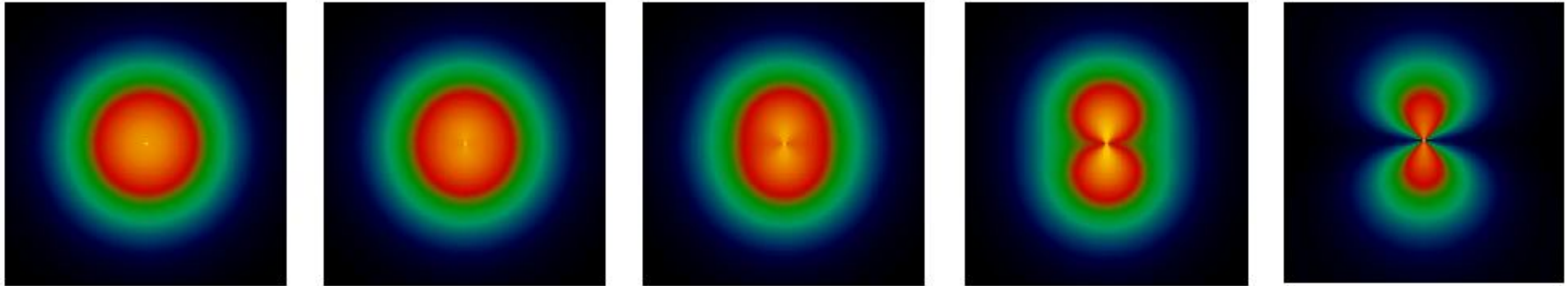
To separate magnetic and nuclear scattering one can use an external magnetic field. In general, nuclear scattering will not depend on the magnetic field (assuming that particles are almost spherical in shape) and magnetic scattering will be strongly field dependent. This follows from the expressions for the corresponding differential cross-sections.

$$\left\langle \frac{d\sigma}{d\Omega} \right\rangle_{nuclear} = (\bar{\rho} - \rho_s)^2 V^2 F^2(Q)$$
$$\frac{d\sigma}{d\Omega}_{magnetic} = const \cdot M^2 V^2 F^2(Q) \sin^2(\Phi)$$

Where V – particle volume, $F(Q)$ – particle form factor, M – particle magnetic moment, Φ - is an angle between particle magnetic moment and the neutron scattering vector Q (in practice, this is the azimuthal angle on the two-dimensional SANS detector).



SANS from magnetic nanoparticles in solution



Simulated 2-dimensional magnetic SANS patterns of superparamagnetic particles in a nonmagnetic matrix. The transition from fully random magnetic orientation of the particles (left figure) to a strong field-parallel alignment (the very right figure) is clearly observed (the hypothetical field is directing horizontally).

SANS from magnetic nanoparticles in solution

Once the magnetic contribution is separated, then by using the contrast variation method it is possible to measure the nuclear radius of the magnetic core as well as the thickness of the polymeric shell.

Obviously we have considered a very simplified situation.

In reality the polydispersity of the particles size, non-spherical shape of the particles, effect of interaction between particles and the very nature of the colloidal systems (which magnetic nanoparticles in solution belong to) being intrinsically non-equilibrium, complicates the data analysis to great extent.



SANS from magnetic nanoparticles in solution

However, it was possible to extract many important characteristics (see e.g. Physics-Uspekhi, 53, 971-993, 2010) like

- Nuclear radius of the magnetic core and thickness of the shell
- Size distribution of the nanoparticles
- Difference between nuclear and magnetic radii of the core
- Temperature behavior of the particles cluster sizes
- Comparison of interparticle interaction for different systems
- etc.

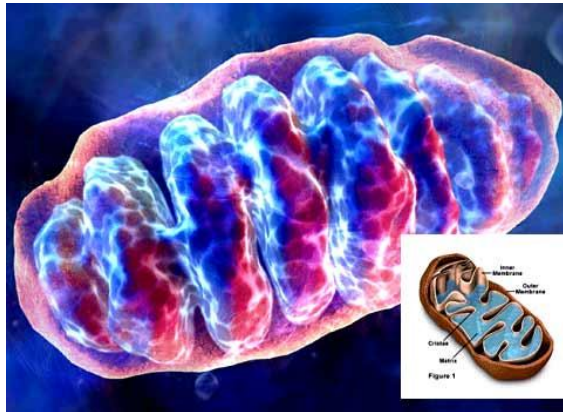


SANS investigations of structural changes in living mitochondria

Inner mitochondrial membrane-

location of energy generating system of mitochondrion

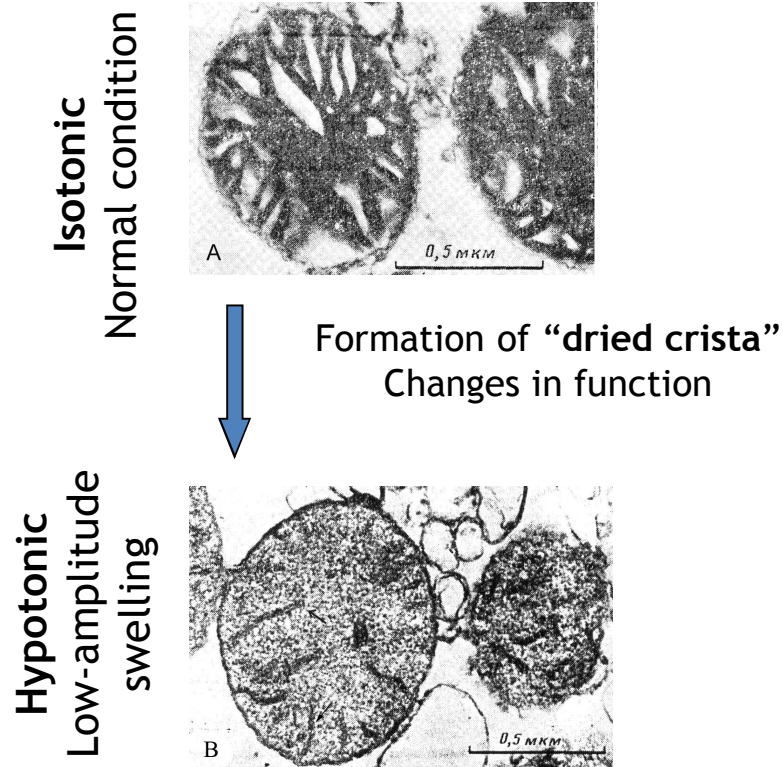




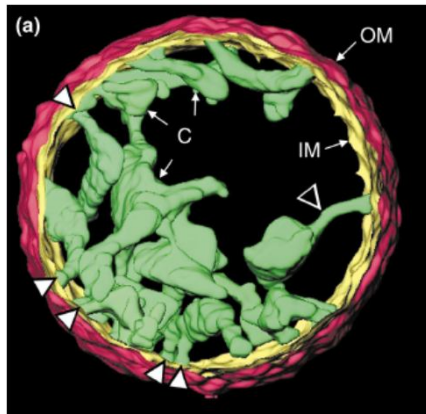
SANS investigations of structural changes in living mitochondria

Hypotonic media induces changes in mitochondrial function:

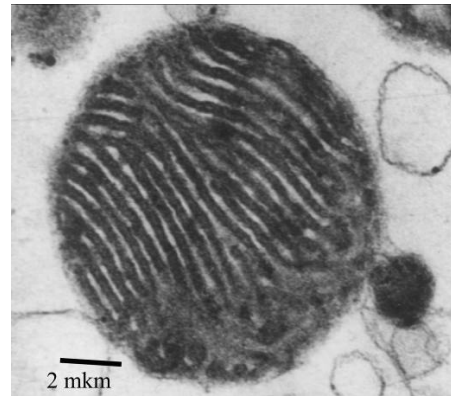
- ❑ regulates the oxidation of fatty acid and other substrates.
- ❑ induces the rising of rate of the electron transport through the respiratory chain
 - ❑ rising of rate and efficiency of ATP synthesis
- ❑ structural changes in polyenzymatic complex of oxidative phosphorylation
 - ❑ switching on the local coupling regime
- ❑ imitation of hormone effect on mitochondria.



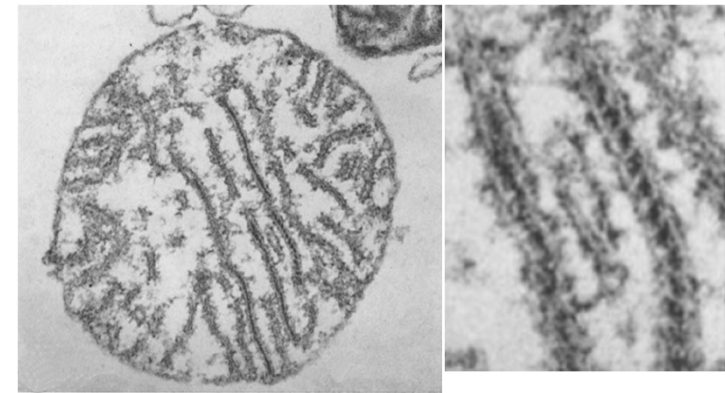
Different packing of inner mitochondrial membranes



Томография митохондрий
печени крысы
*T.G. Frey, TIBS V.25, 2000, p.
319-324*

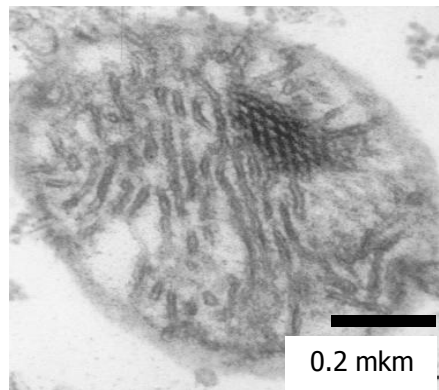


Митохондрии сердца крысы



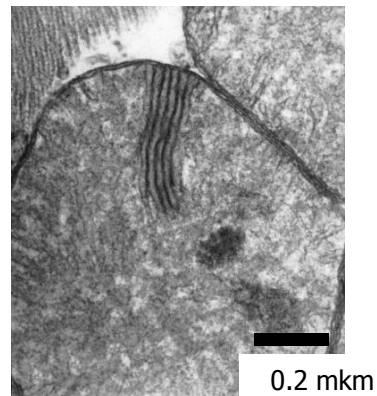
Предшественники гексагональной фазы
Митохондрии сердца быка

J. Hall and F. Crane, The J. Cell Biology, V.48, 1971, p. 420-425

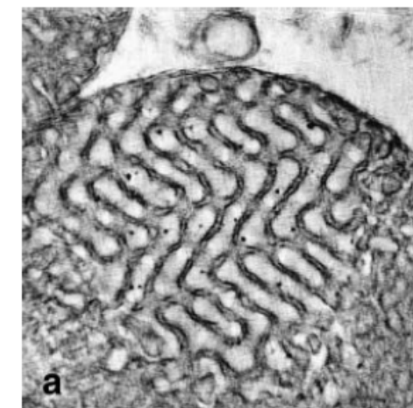


Гексагональная фаза

Митохондрии сердца крысы при апоптозе
Сапрунова В.Б., Бакеева Л.Е., Ягужинский Л.С. // Цитология.
2003. Т. 45. С. 1073-1082

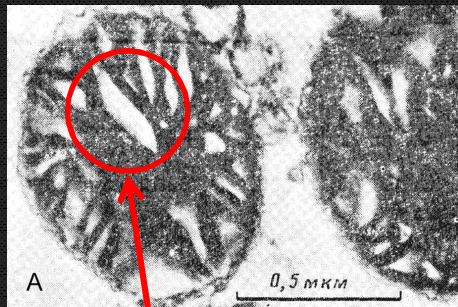
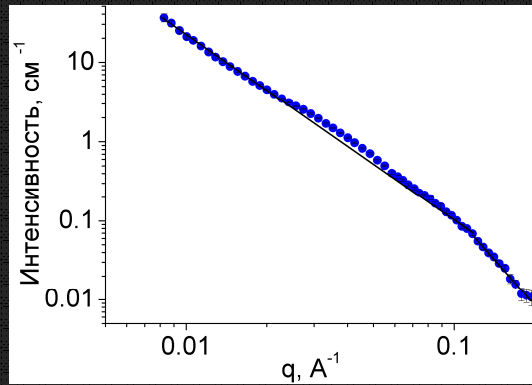


0.2 mkm



Кубическая фаза

Митохондрии голодающей амёбы *C. Carolinensis*.
Окислительный стресс
Y. Deng et al, Protoplasma (2002) 219: 160-167

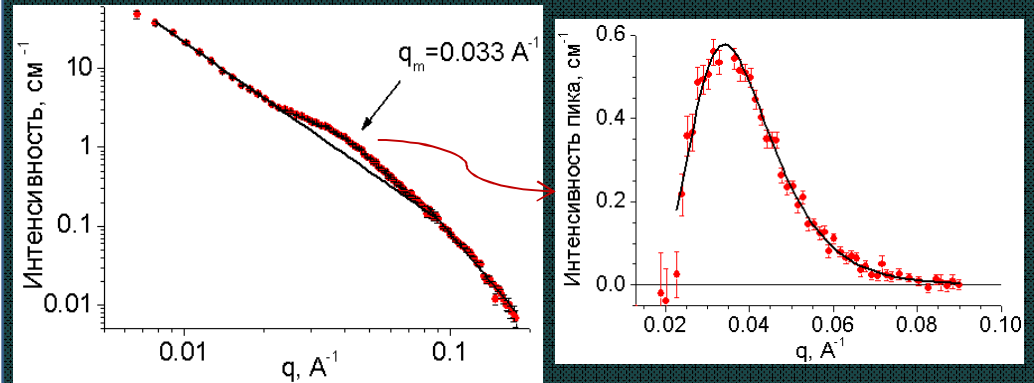


Expanded disordered cristae

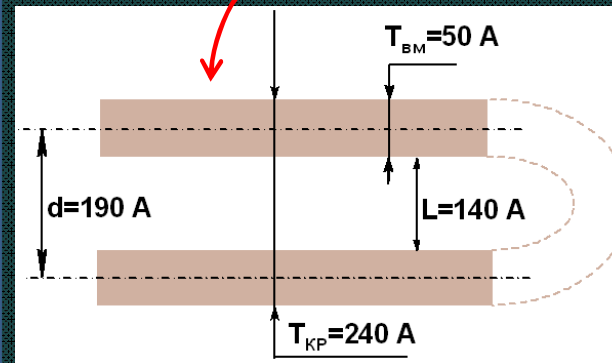
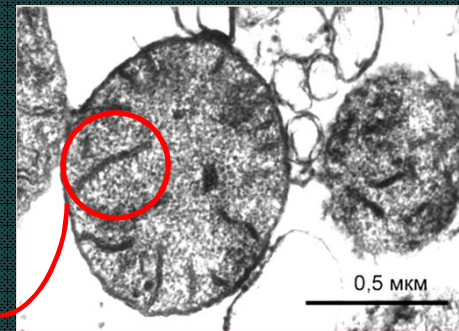
No peaks

Krasinskaya I. et al. // *Biokhimiya* 1989, V. 54: 1550-1556.
 Murugova T. et al. // *Biophysics* 2006, V.51: 882-886.

Hypotonic conditions

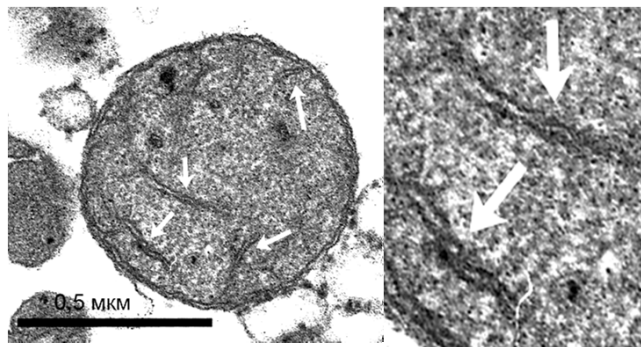


Extracted interference peak

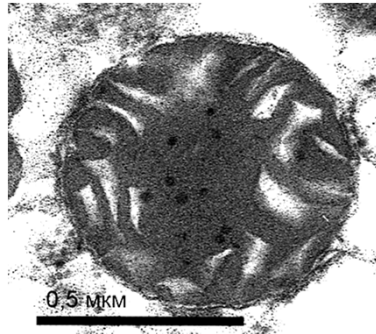


Scheme of dried ordered cristae formed under hypotonic conditions

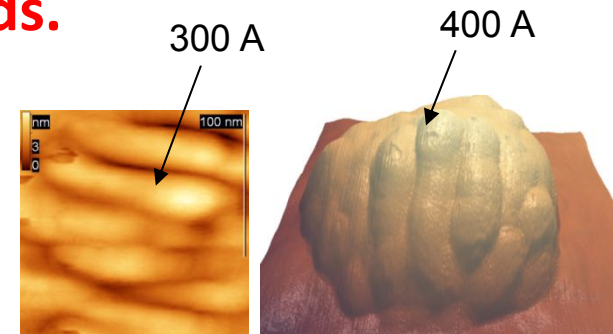
Comparison of double membrane structural parameters obtained by different methods.



Hypotonic medium

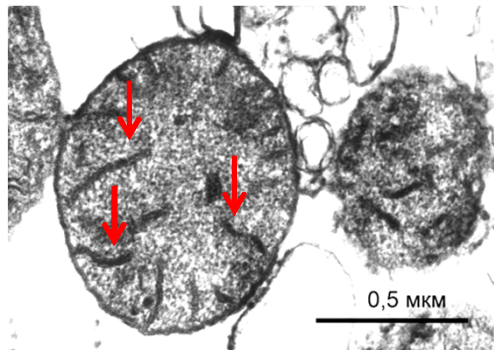


Isotonic medium



AFM data

Electron microscopy



Красинская И.П., Литвинов И.С., Захаров С.Д.,
Бакеева Л.Е., Ягужинский Л.С. (1989), *Биохимия*,
54(9), 1550-1556.

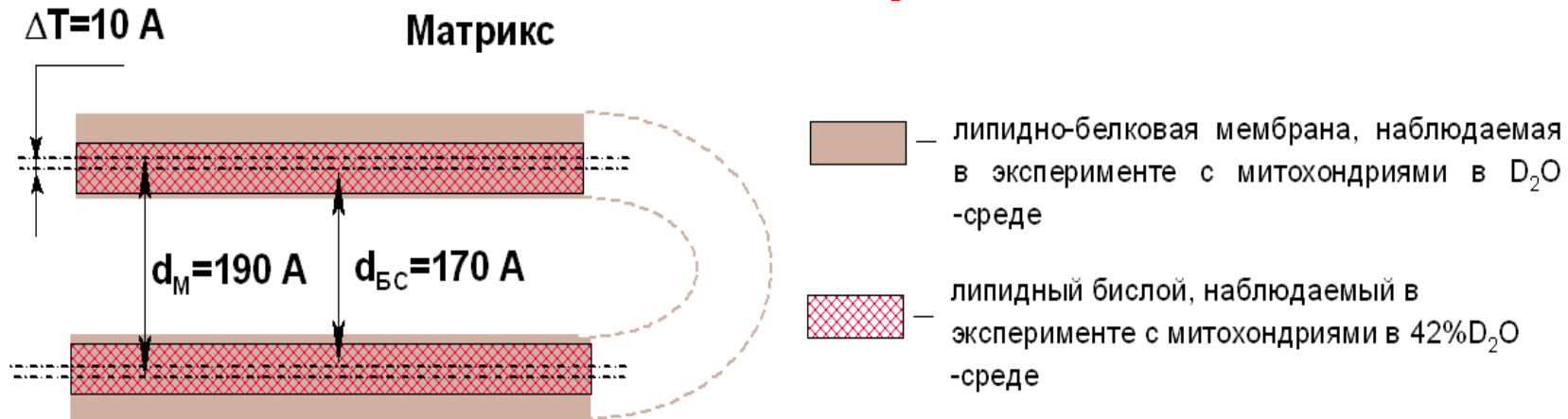
Method	Double membrane structure thickness, Å
SANS (alive, statistical average)	240±10 *265±10
Electron microscopy (dead, individual)	220±20 #250±30
AFM (dead, individual)	300-400

* С учетом литературных данных о толщине внутренней митохондриальной мембраны
Фотография из работы Красинская И.П., Литвинов И.С., Захаров С.Д., Бакеева Л.Е.,
Ягужинский Л.С. (1989), *Биохимия*, 54(9), 1550-1556.

Dubrovin E.V., Murugova T.N., Motovilov K.A., Yaguzhinskii L.S., Yaminskii I.V.
Nanotechnologies in Russia 2009;4:876-880.



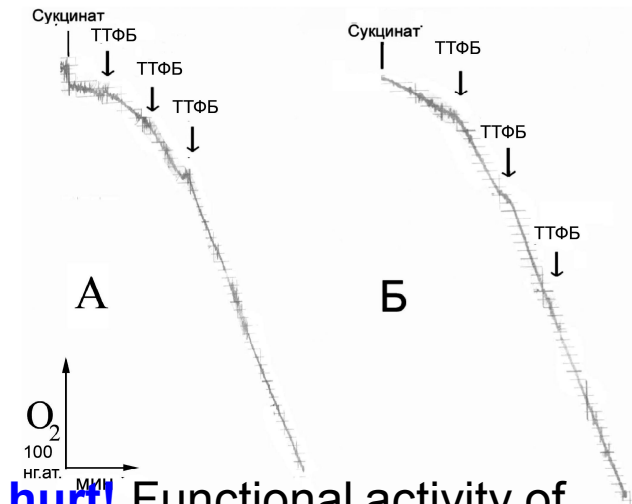
Observation of asymmetric distribution of proteins with respect to the lipid membranes using contrast variation technique.



Distance between centres of double layer lipid membranes $d_{BC} = 170 \text{ \AA}$

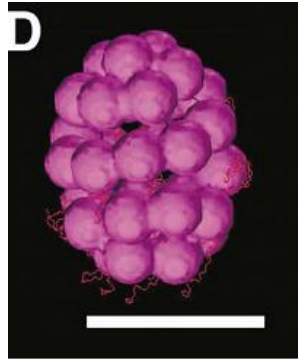
Distance between centres of double layer lipid membranes with imbedded proteins $d_M = 190 \text{ \AA}$.

Centre of gravity for proteins is shifted towards matrix.

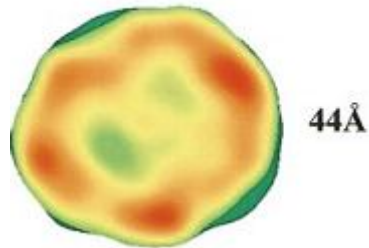


Neutrons do not hurt! Functional activity of mitochondria before (A) and after (B) SANS experiment

Structure of the α -crystalline by small-angle neutron scattering (FLNP in collaboration with the Institute of biochemical physics RAN)



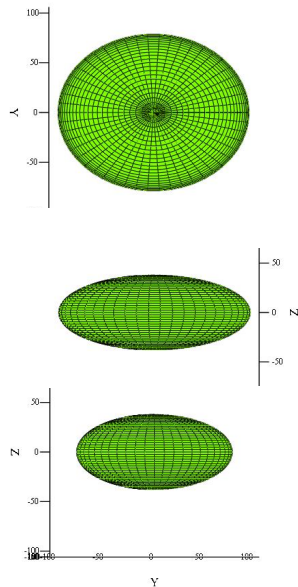
A model of α B-crystallin with bound α -lactalbumin. **Electron microscopy**. Scale bar represents 100 Å. J. Horwitz, *Exp. Eye Res.* (2003) 76, 145-153



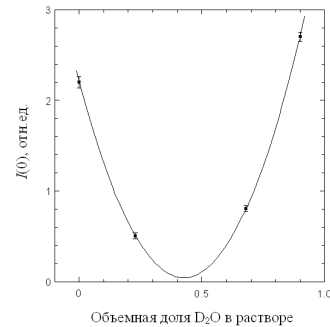
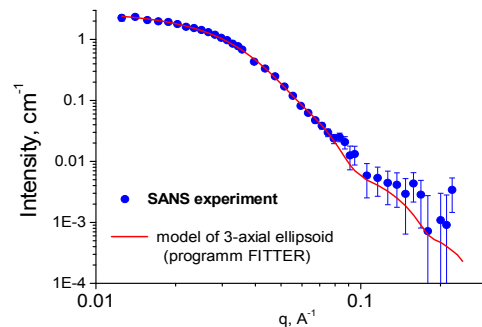
Single particle reconstructions of native α -crystallin from **Electron microscopy**. D.A. Haley et al., *J. Mol. Biol.* (2000) 298, 261±272

- Oligomer protein (20-40 monomers, mass 400-1000 kDa)
 - Polydisperse system
 - Dynamical system (exchange of monomers)
 - Difficult to crystallize
- This is the main component of the eye lens of vertebrates (200 mg/ml), which ensures a correct refractive index
- Exhibits chaperone like activity (forms soluble complexes with the destabilized proteins and prevents their non-specific aggregation and uncontrolled denaturation)
- Slows down the age related opacification of the eye lens (cataract)

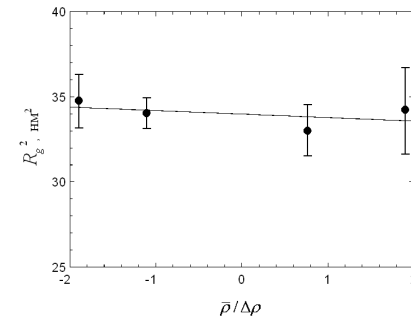
Structure of the α -crystalline by small-angle neutron scattering (FLNP in collaboration with the Institute of biochemical physics RAN)



Approximation of the SANS data by the model of three-axis ellipsoid.



Experimental dependence of the scattering intensity at zero angle versus volume fraction of D_2O in the buffer solution and its approximation by parabolic function.



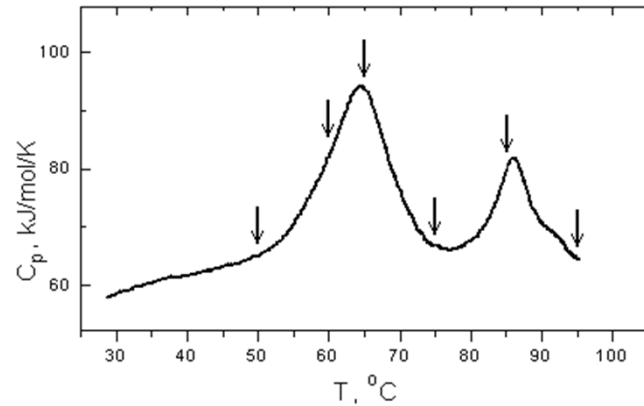
Experimental dependence of V crystalline radius of gyration versus contrast in buffer solution with different D_2O concentrations.

- Uniform distribution of the scattering density in the protein parts inaccessible for water penetration.
- Absence of voids into which water cannot penetrate.
- All sub-units are equally accessible for water (perhaps for other low molecular weight substances, including bio active compounds).

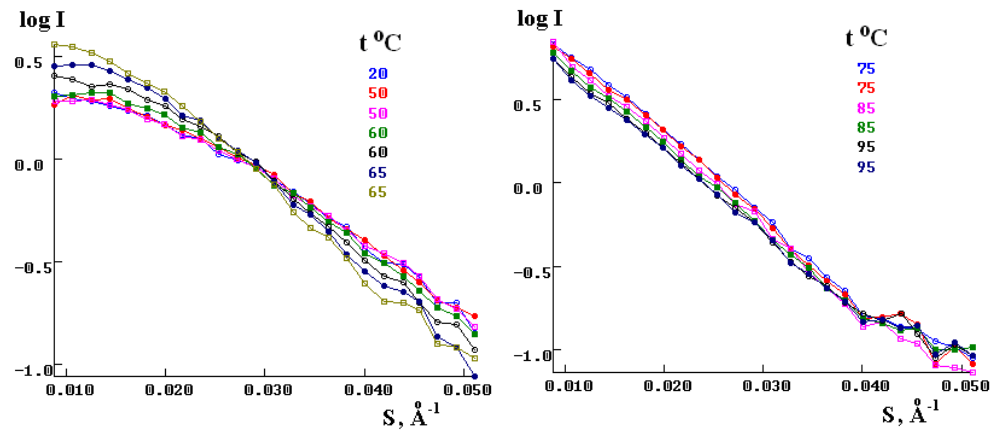
А. В. Кривандин, Т. Н. Муругова, А. И. Куклин, К. О. Муранов, Н. Б. Полянский, В. Л. Аксенов, М. А. Островский, Биохимия **75** (11), 1499-1507 (2010).



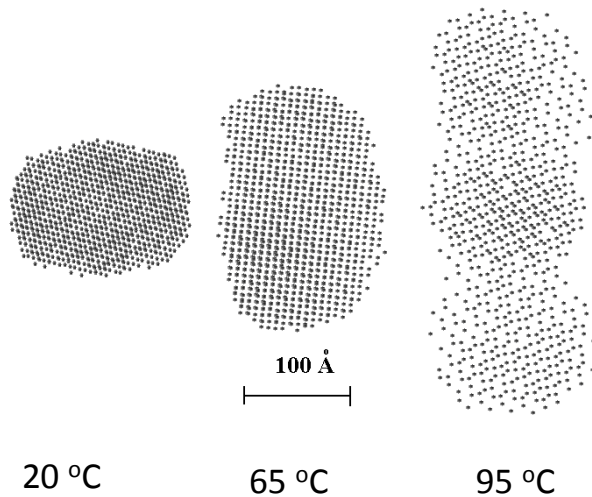
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DSC data. Arrows indicate temperatures at which SANS data were measured.



Changes of SANS curves for α crystalline in 90% D_2O -buffer during heating from 20 to 95°C.



- Protein dimerization at 60-65 °C
- Rod-like form at 75-95 °C

A. V. Krivandin, A. I. Kuklin, K. O. Muranov, **T. N. Murugova**, S. S. Kozlov, N. K. Genkina, "Heat-induced structural transitions of alpha-crystallin studied by small-angle neutron scattering", Journal of Physics: Conference Series, 2012, Vol. 351, 012008.

Complex study of nanostructures

- Electron microscopy
- Dynamic light scattering
- Static light scattering
- Quasi elastic neutron scattering
- X-ray and neutron diffraction
- X-ray and neutron reflectometry
- Small-angle X-ray and neutron scattering
- Sedimentation analysis
- Spectrometry
- Photometry
- Calorimetry

