

Физика на елементарните частици

Lecture 1
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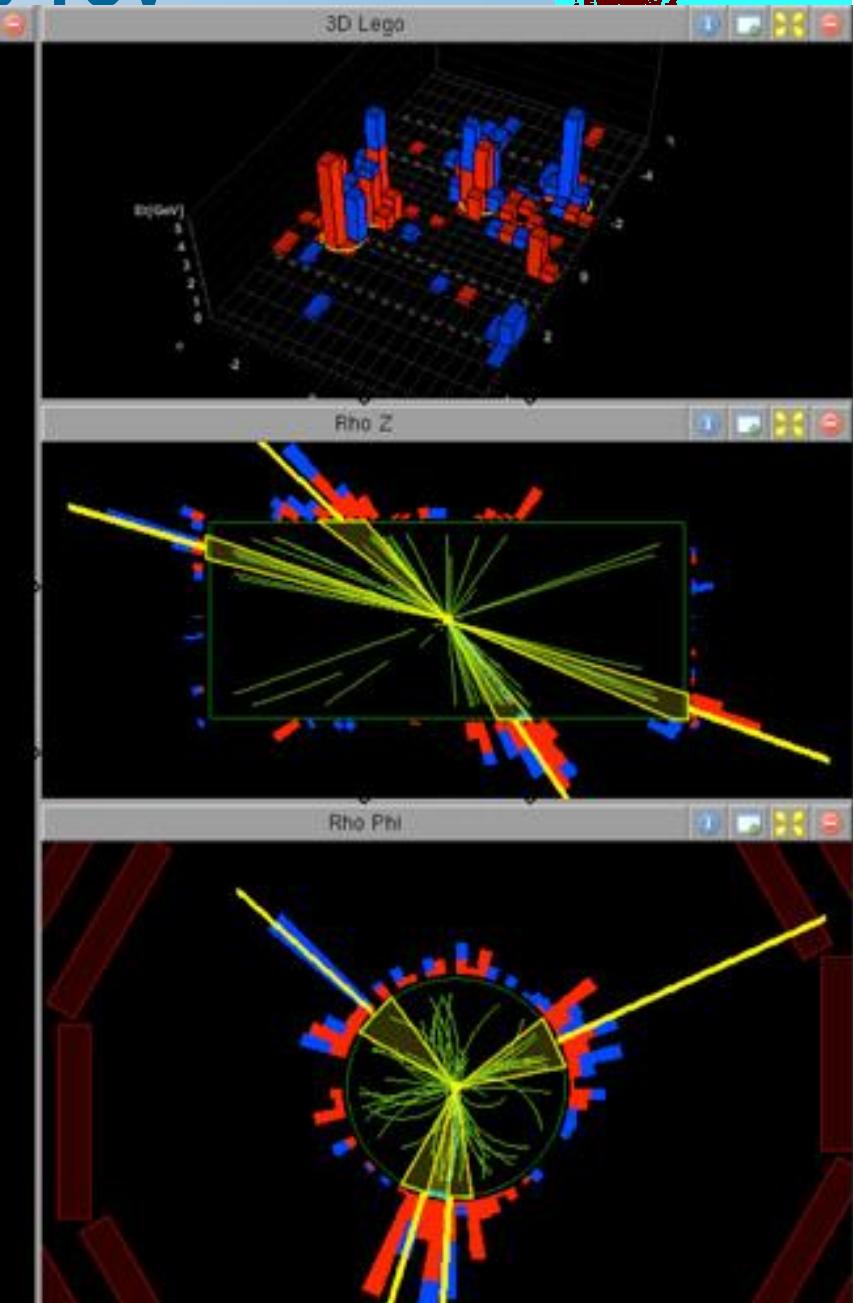
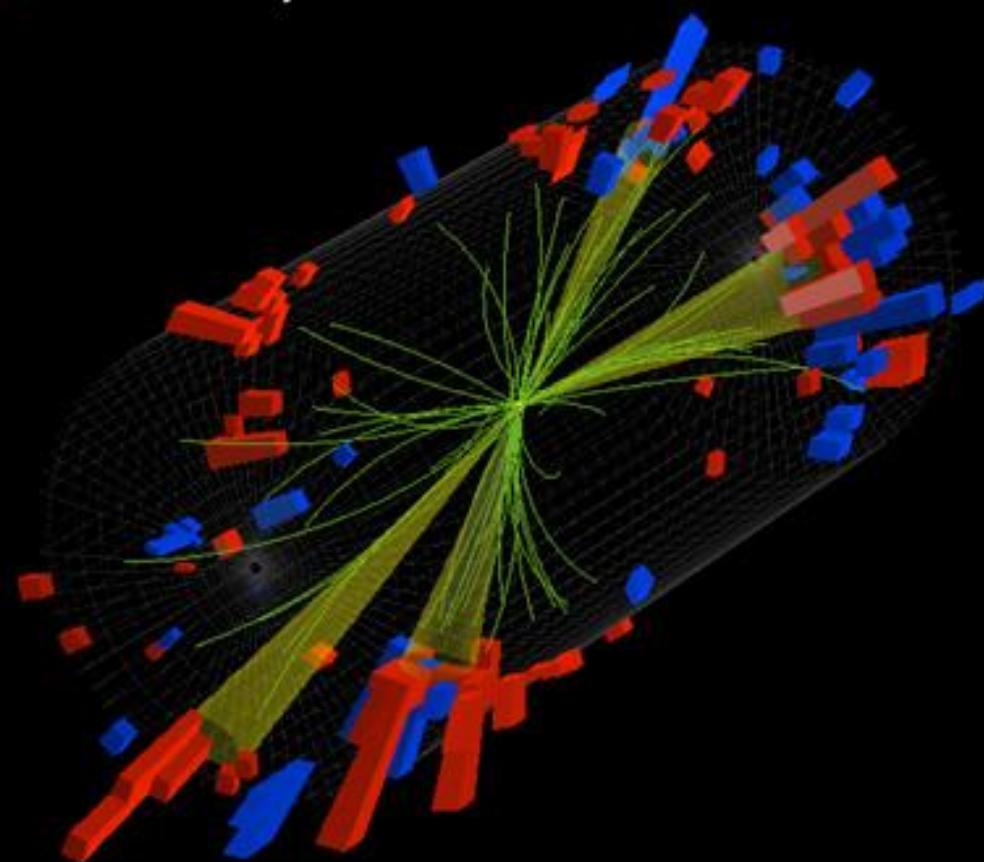
May 2010, Sofia



Collisions at 2.36 TeV

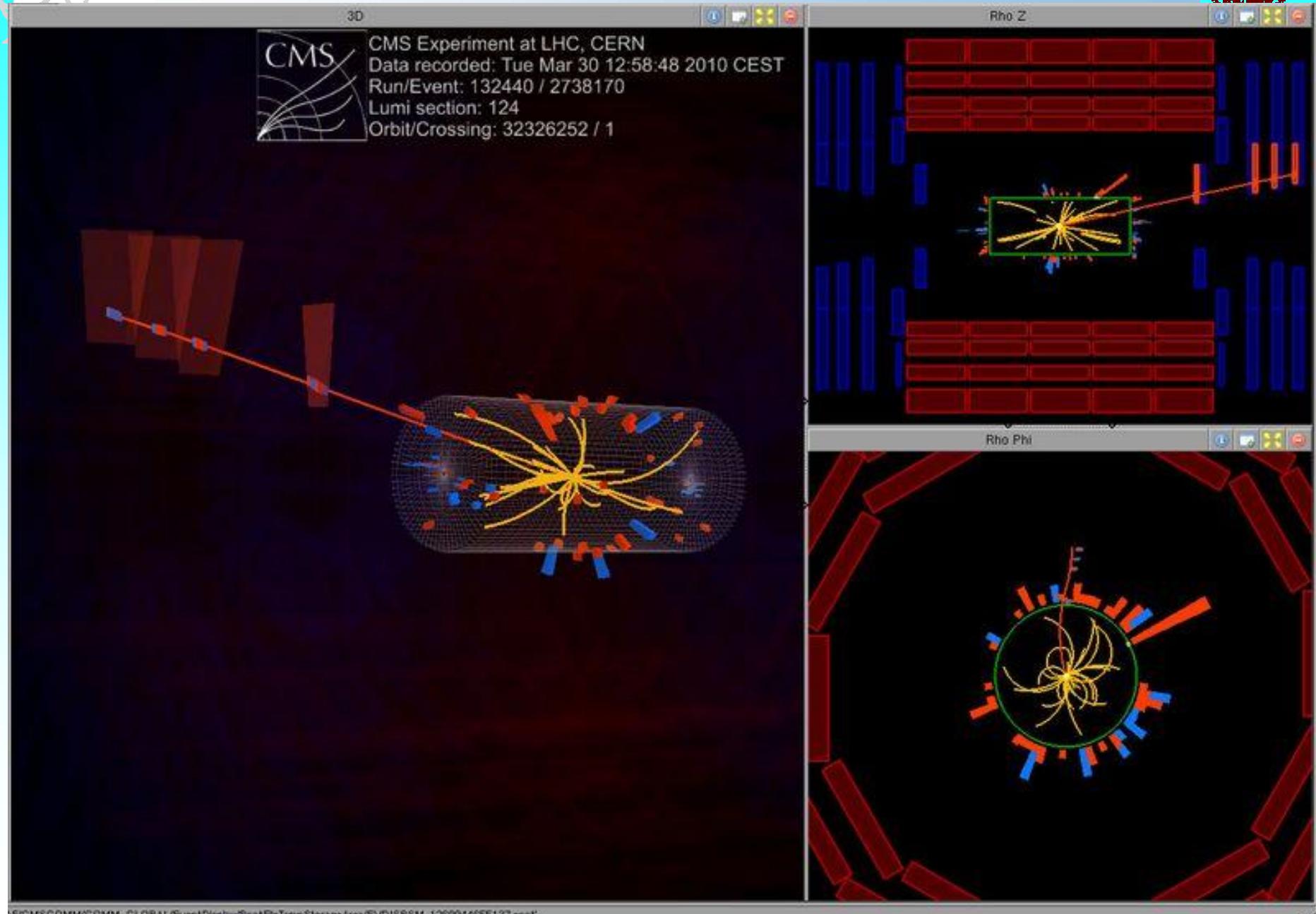


CMS Experiment at the LHC, CERN
Date Recorded: 2009-12-14
Run/Event: 124120/6613074
Candidate Multijet Event at 2.36 TeV





Collisions at 7 TeV





Collisions at 7 TeV



CMS Experiment at the LHC, CERN

Data recorded: 2010-Mar-30 11:04:14.111090 GMT(13:04:14 CEST)

Run: 132440
Event: 3087931
Lumi section: 138
Orbit: 35985009
Crossing: 1





Що е то?



Физика на елементарните частици



Опитва се да отговори на два фундаментални въпроса

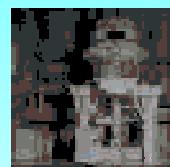
-Кои са елементарните съставящи на материята?

-Кои са фундаменталните сили контролиращи
тяхното поведение ?

Instruments



Accelerators
LHC, LEP



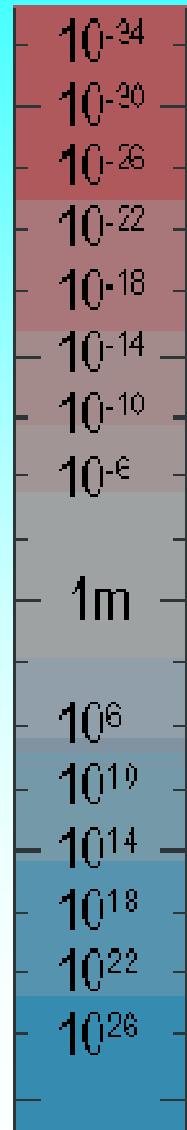
(Particle beams)
Electron
Microscope
Microscope



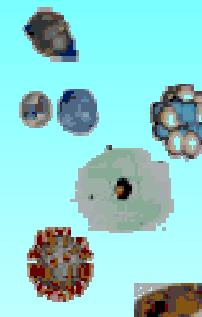
Telescope



Radio
Telescope



Observables



SJSY particle?
Higgs? (range of
Z,W
Proton
Nuclei
Atom
Virus
Cell



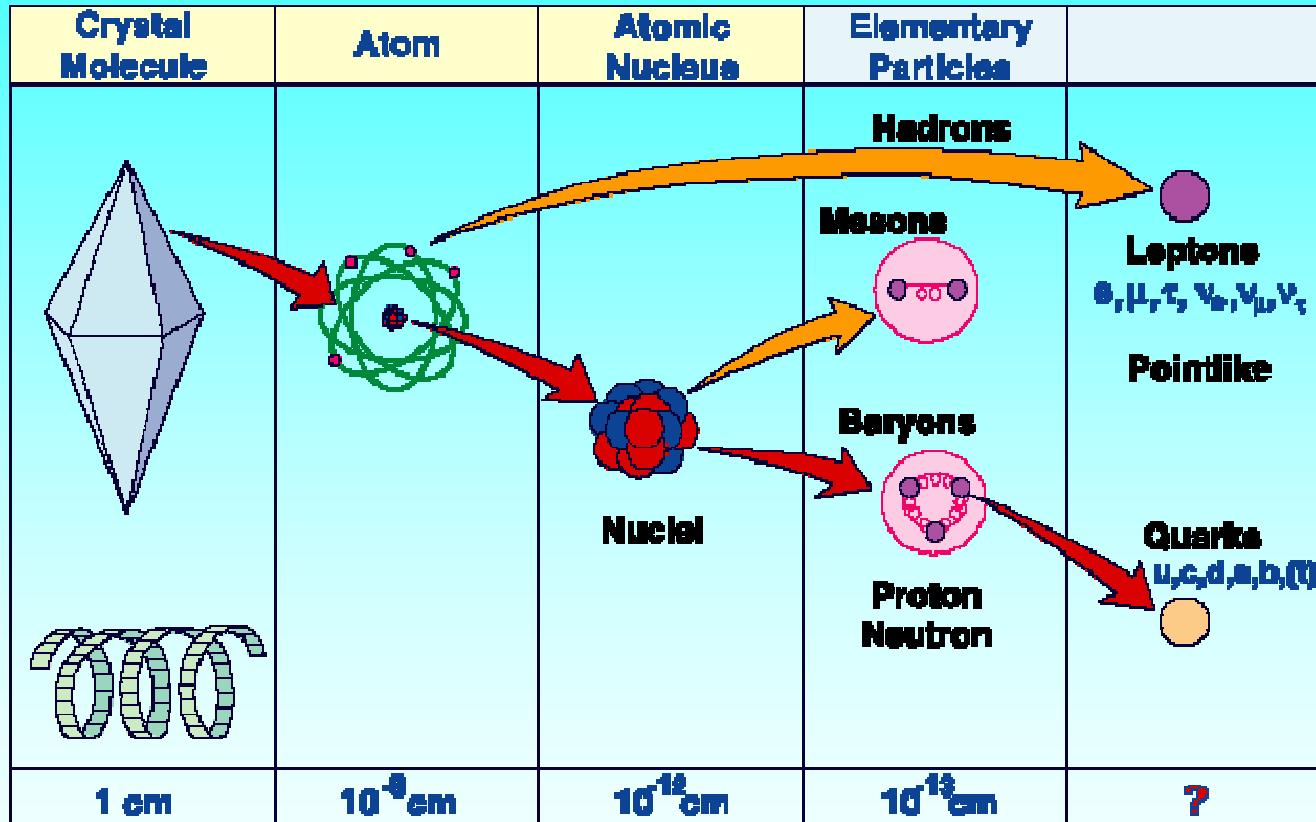
Earth radius
Earth to Sun



Galaxies
Radius of observable
Universe



Структура на материята

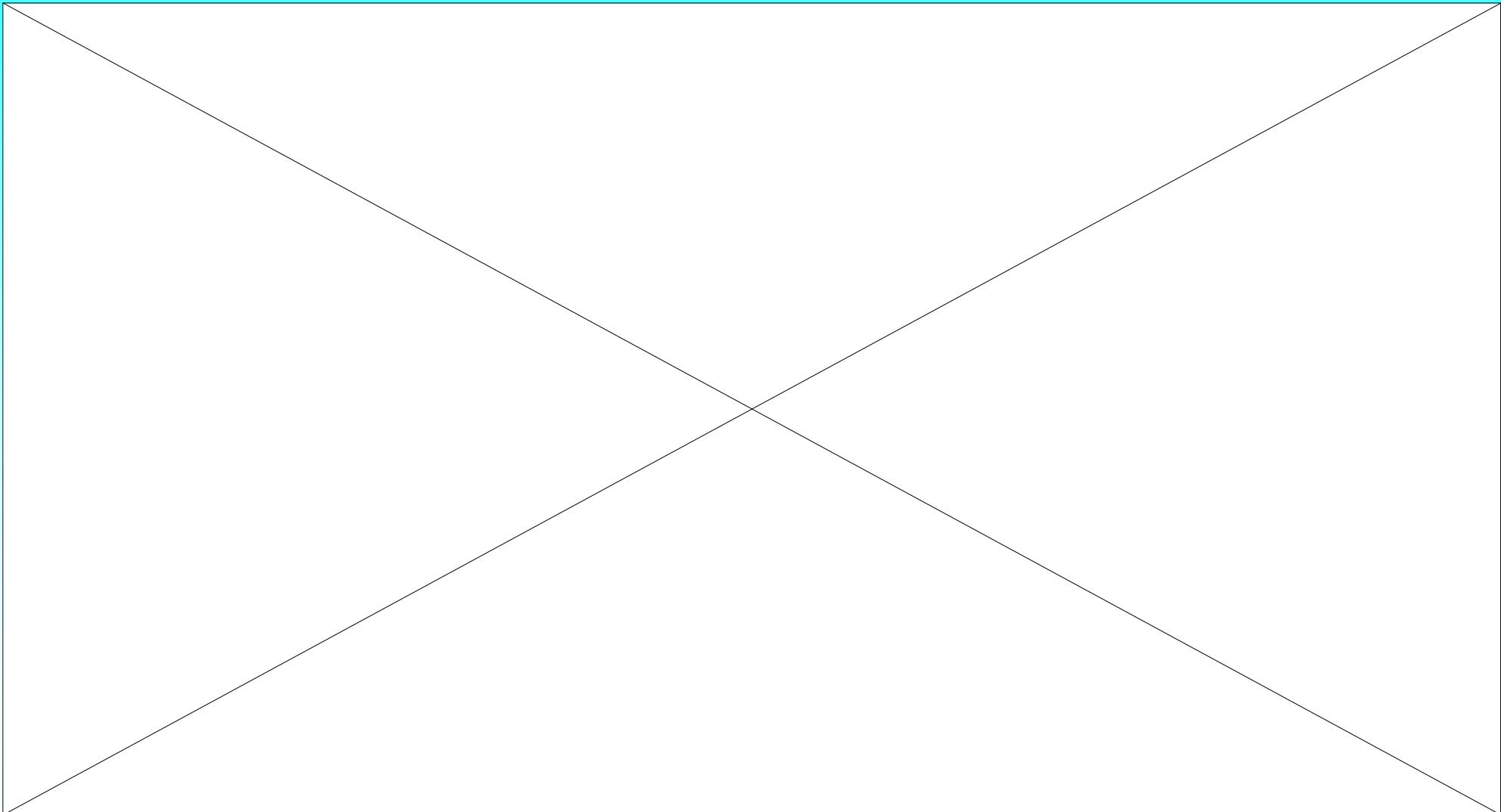


$\approx 10^4$

Физика на елементарните частици



Структура на материјата





Единици във физиката на елементарните частици

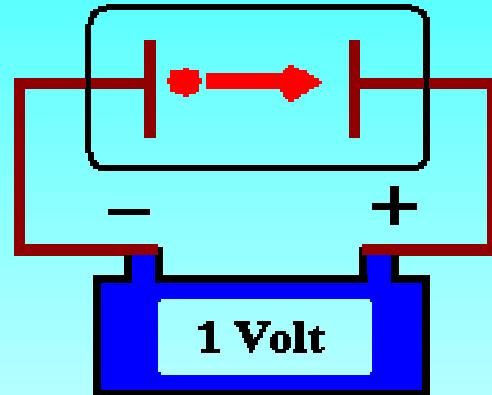


Energy

1 electron-Volt (eV):

the energy of a particle with electric charge = $|e|$,
initially at rest, after acceleration by a difference
of electrostatic potential = 1 Volt

$$(e = 1.60 \times 10^{-19} \text{ C})$$



$$1 \text{ eV} = 1.60 \times 10^{-19} \text{ J}$$

Multiples:

$$\begin{aligned} 1 \text{ keV} &= 10^3 \text{ eV} ; & 1 \text{ MeV} &= 10^6 \text{ eV} \\ 1 \text{ GeV} &= 10^9 \text{ eV}; & 1 \text{ TeV} &= 10^{12} \text{ eV} \end{aligned}$$

Energy of a proton in the LHC :

$$7 \text{ TeV} = 1.12 \times 10^{-6} \text{ J}$$

(the same energy of a body of mass = 1 mg moving at speed = 1.5 m /s)



Маса, разстояние, енергия, температура



These are related quantities

In particle physics the unit of energy is the **electron volt**.

1 electron volt (eV) = Energy gained by an electron in passing through a voltage difference of 1 V

$$E=mc^2$$

$$c=3 \cdot 10^8 \text{ m/s}$$

speed of light

$$E=kT$$

$$k=10^{-4} \text{ eV K}^{-1}$$

Boltzmann's constant

$$E=h\nu/\lambda$$

$$h=4 \cdot 10^{-15} \text{ eV s}$$

Planck's constant

Mass of electron

0.5 million eV (MeV)

Mass of proton

1 Giga eV (GeV)

1 eV \sim 10,000 K

1 GeV \sim 1 femtometre (fm) = 10^{-15} m

General Relativity depends on c and G (Newton's constant), QM depends on \hbar . Natural unit of length is given by is called Planck length $\sim 10^{-35} \text{ m}$

$$\sqrt{\hbar G / c^3}$$



Фундаментални съставящи



1894 – 1897: Откритие на електрона



Study of “cathode rays”: electric current in tubes at very low gas pressure (“glow discharge”)

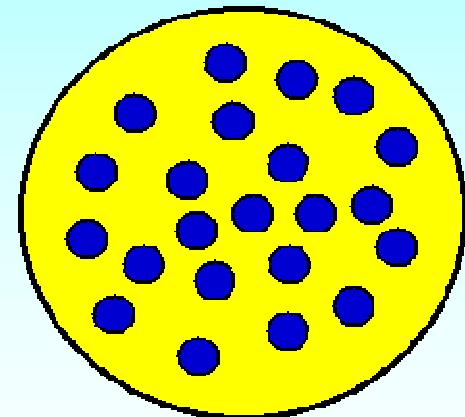
Measurement of the electron mass: $m_e \gg M_H/1836$

“Could anything at first sight seem more impractical than a body which is so small that its mass is an insignificant fraction of the mass of an atom of hydrogen?” (J.J. Thomson)

→ ATOMS ARE NOT ELEMENTARY

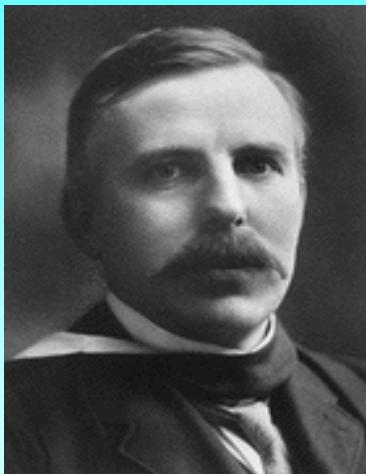
Thomson's atomic model:

- § Electrically charged sphere
- § Radius $\sim 10^{-8}$ cm
- § Positive electric charge
- § Electrons with negative electric charge embedded in the sphere

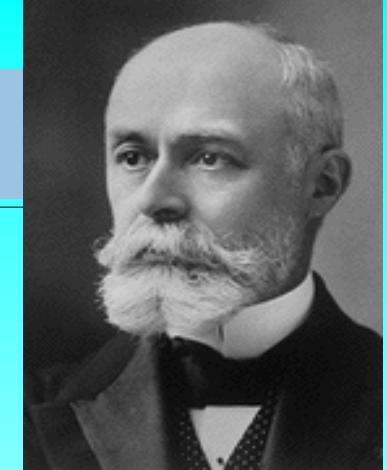




1896: Естествена радиоактивност

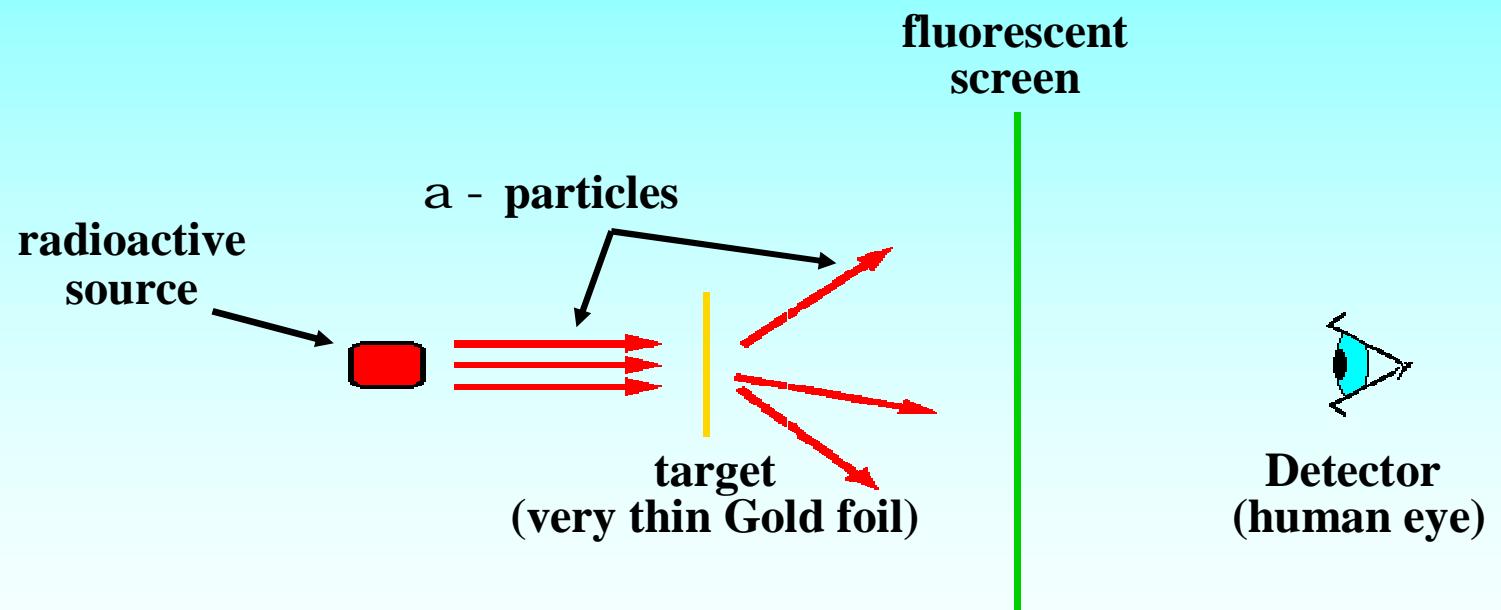


Ernest Rutherford



Henri Becquerel

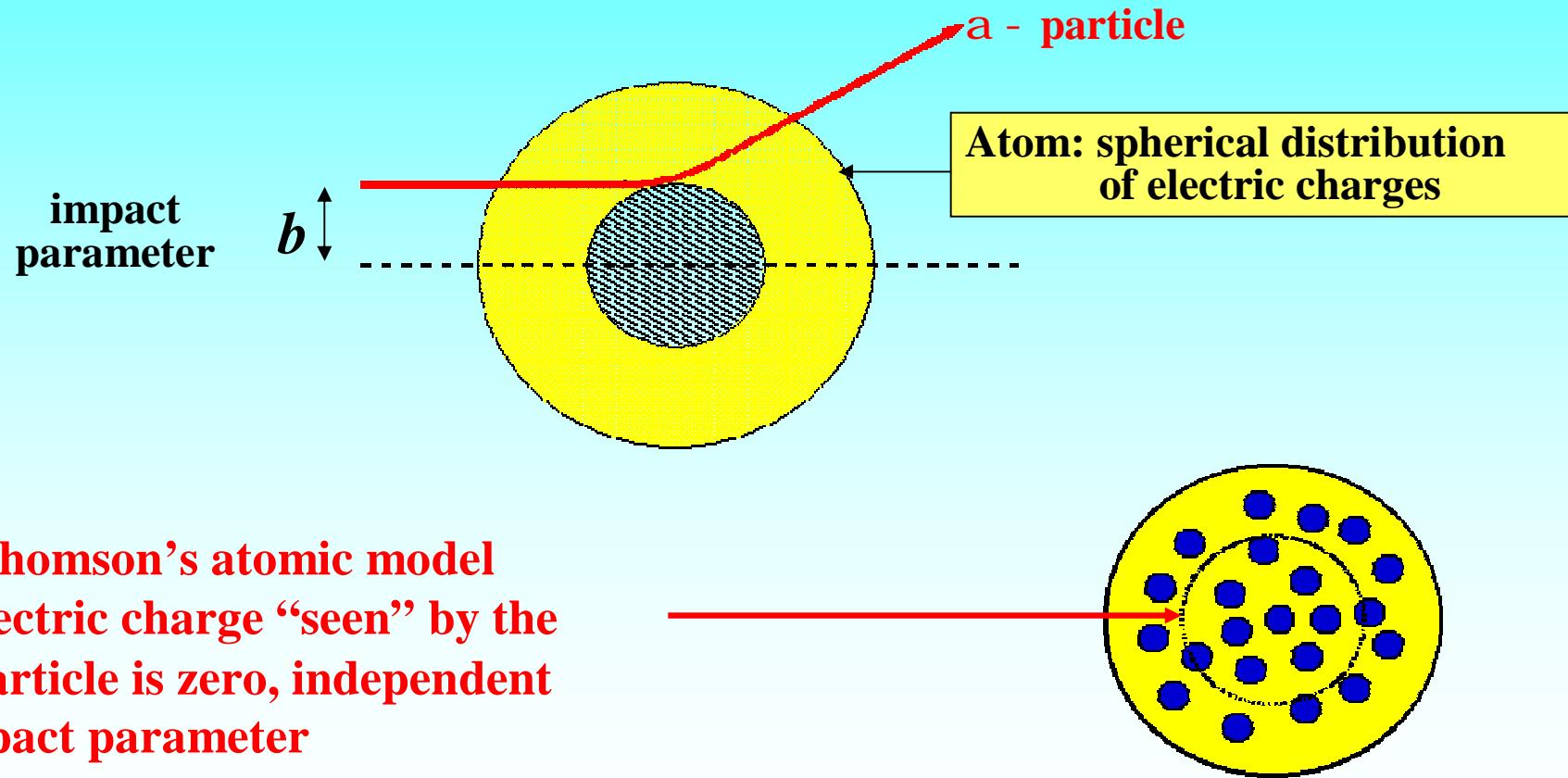
1909 - 13: Rutherford's scattering experiments
Discovery of the atomic nucleus



a - particles: nuclei of Helium atoms spontaneously emitted by heavy radioactive isotopes

Typical a – particle velocity $\gg 0.05 c$ (c : speed of light)

a – atom scattering at low energies is dominated by Coulomb interaction

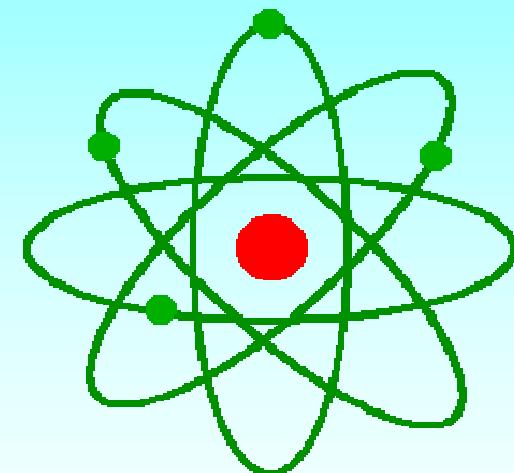


For Thomson's atomic model
the electric charge “seen” by the
a – particle is zero, independent
of impact parameter

Þ no significant scattering at large angles is expected

significant scattering of α – particles at large angles, consistent with scattering expected for a sphere of radius \approx few $\times 10^{-13}$ cm and electric charge = Ze , with $Z = 79$ (atomic number of gold) and $e = |\text{charge of the electron}|$

→ **an atom consists of
a positively charged nucleus
surrounded by a cloud of electrons**



Nuclear radius $\gg 10^{-13}$ cm $\gg 10^{-5} \times$ atomic radius

Mass of the nucleus \gg mass of the atom
(to a fraction of 1%)



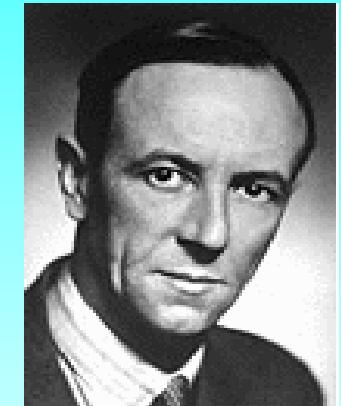
Откритие на неутрона



**Neutron: a particle with mass » proton mass
but with zero electric charge (Chadwick, 1932)**

Solution to the nuclear structure problem:

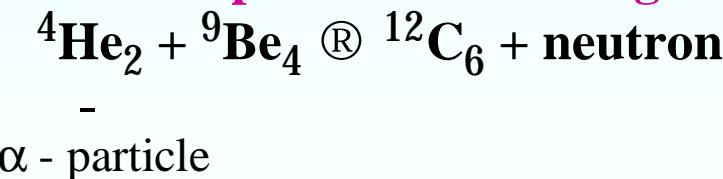
**Nucleus with atomic number Z and mass number A :
a bound system of Z protons and $(A - Z)$ neutrons**



Nitrogen anomaly: no problem if neutron spin = $\frac{1}{2}\hbar$

Nitrogen nucleus ($A = 14$, $Z = 7$): 7 protons, 7 neutrons = 14 spin $\frac{1}{2}$ particles
↳ total spin has integer value

**Neutron source in Chadwick's experiments: a ^{210}Po radioactive source
(5 MeV α -particles) mixed with Beryllium powder ↳ emission of
electrically neutral radiation capable of traversing several centimetres of Pb:**





Принцип на Паули



In Quantum Mechanics the electron orbits around the nucleus are “quantized”: **only some specific orbits (characterized by integer quantum numbers) are possible.**

$$R_n = \frac{4\pi\varepsilon_0 \hbar^2 n^2}{me^2} \approx 0.53 \times 10^{-10} n^2 \text{ [m]}$$

$$\left. \begin{array}{l} m = m_e m_p / (m_e + m_p) \\ n = 1, 2, \dots \end{array} \right\}$$

$$E_n = -\frac{me^4}{2(4\pi\varepsilon_0)^2 \hbar^2 n^2} \approx -\frac{13.6}{n^2} \text{ [eV]}$$

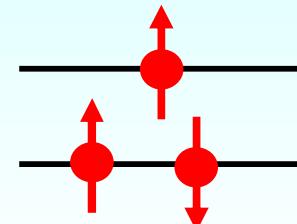
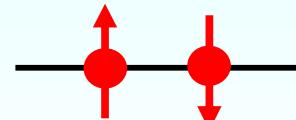
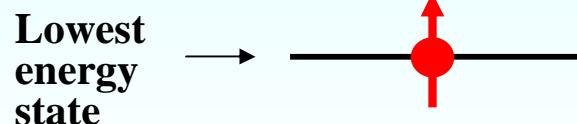
In atoms with $Z > 2$ only two electrons are found in the innermost orbit – WHY?

ANSWER (Pauli, 1925): two electrons (spin = $\frac{1}{2}$) can never be in the same physical state

Hydrogen ($Z = 1$)

Helium ($Z = 2$)

Lithium ($Z = 3$)



Wolfgang Pauli

Pauli's exclusion principle applies to all particles with half-integer spin (collectively named Fermions)

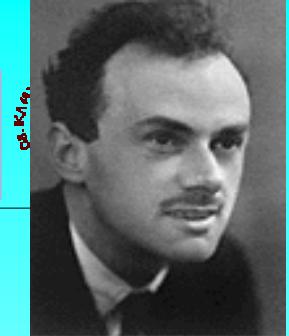
L. Litov

Физика на елементарните частици

Sofia, April 2009



Антиматерия



Discovered “theoretically” by P.A.M. Dirac (1928)

Dirac’s equation: a relativistic wave equation for the electron

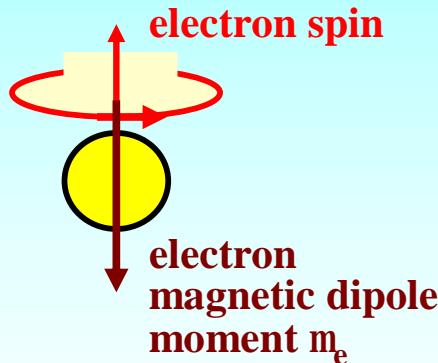
P.A.M. Dirac

Two surprising results:

§ Motion of an electron in an electromagnetic field:

presence of a term describing (for slow electrons) the potential energy of a magnetic dipole moment in a magnetic field

► existence of an intrinsic electron magnetic dipole moment opposite to spin



$$\mu_e = \frac{e\mathbf{h}}{2m_e} \approx 5.79 \times 10^{-5} \text{ [eV/T]}$$

§ For each solution of Dirac’s equation with electron energy $E > 0$ there is another solution with $E < 0$

What is the physical meaning of these “negative energy” solutions ?

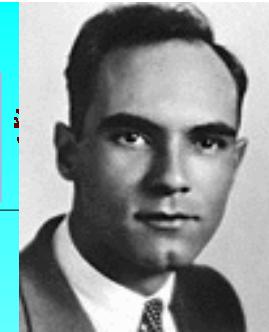
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Физика на елементарните частици

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Експериментално наблюдение на антиматерия



(C.D. Anderson, 1932)

Measure particle momentum and sign of electric charge from magnetic curvature

Detector: a Wilson cloud – chamber (visual detector based on a gas volume containing vapour close to saturation) in a magnetic field, exposed to cosmic rays

→ projection of the particle trajectory in a plane perpendicular to \vec{B} is a circle

$$\text{Lorentz force} \quad \vec{f} = e\vec{v} \times \vec{B}$$

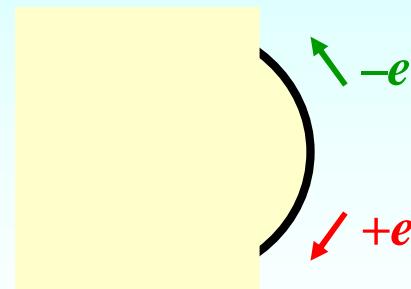
Circle radius for electric charge $|e|$:

p_{\perp} : momentum component perpendicular to magnetic field direction

$$R [\text{m}] = \frac{10 p_{\perp} [\text{GeV}/c]}{3B [\text{T}]}$$

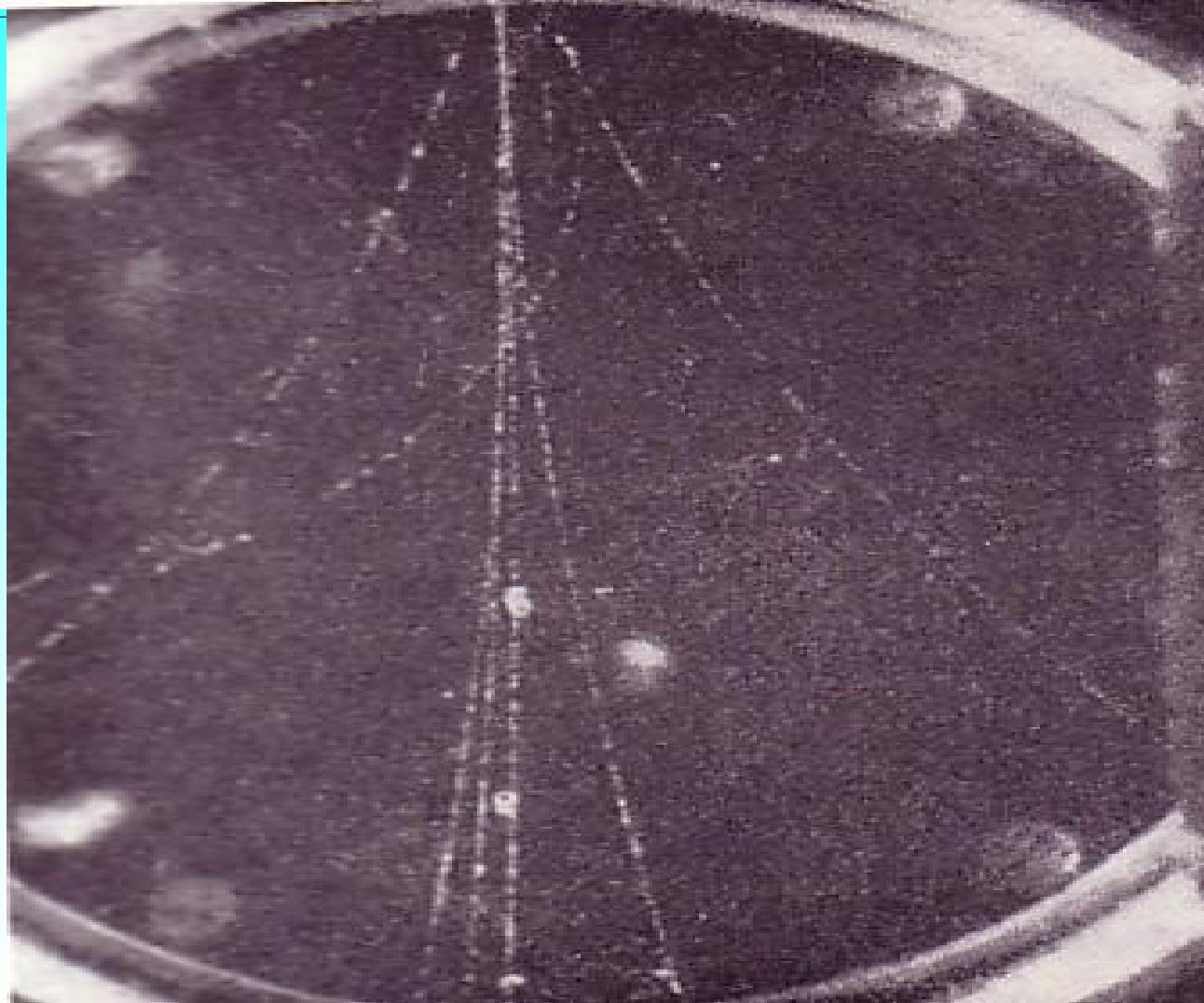
NOTE: impossible to distinguish between positively and negatively charged particles going in opposite directions

⇒ need an independent determination of the particle direction of motion





Experimental confirmation of antimatter



Cosmic-ray “shower”
containing several $e^+ e^-$ pairs

Физика на элементарните частици

L. Litov

Sofia, April 2009



Неутрино



December 1930: public letter sent by W. Pauli to a physics meeting in Tübingen

Zürich, Dec. 4, 1930

Dear Radioactive Ladies and Gentlemen,

...because of the “wrong” statistics of the N and ${}^6\text{Li}$ nuclei and the continuous β -spectrum, I have hit upon a desperate remedy to save the law of conservation of energy. Namely, the possibility that there could exist in the nuclei electrically neutral particles, that I wish to call neutrons, which have spin $1/2$ and obey the exclusion principle The mass of the neutrons should be of the same order of magnitude as the electron mass and in any event not larger than 0.01 proton masses. The continuous β -spectrum would then become understandable by the assumption that in β -decay a neutron is emitted in addition to the electron such that the sum of the energies of the neutron and electron is constant.

..... For the moment, however, I do not dare to publish anything on this idea

So, dear Radioactives, examine and judge it. Unfortunately I cannot appear in Tübingen personally, since I am indispensable here in Zürich because of a ball on the night of 6/7 December.

W. Pauli



Теория на β–разпада (V-A модел)



Enrico Fermi

b^- decay: $n \xrightarrow{\beta^-} p + e^- + \bar{n}$

b^+ decay: $p \xrightarrow{\beta^+} n + e^+ + n$ (e.g., $^{14}\text{O}_8 \rightarrow ^{14}\text{N}_7 + e^+ + \nu$)

n : **the particle proposed by Pauli**
(named “neutrino” by Fermi)

\bar{n} : **its antiparticle (antineutrino)**

(E. Fermi, 1932-33)

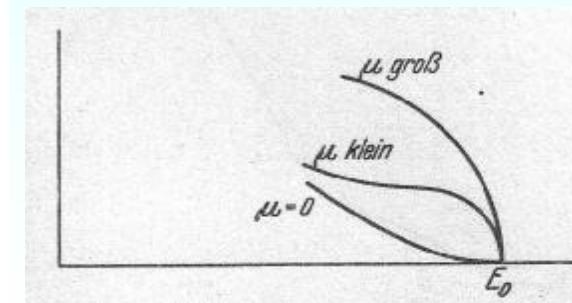
Fermi’s theory: a point interaction among four spin $\frac{1}{2}$ particles, using
the mathematical formalism of creation and annihilation
operators invented by Jordan
 \bar{p} particles emitted in b^- decay need not exist before emission –
they are “created” at the instant of decay

Prediction of b^- decay rates and electron energy spectra as a function of
only one parameter: Fermi coupling constant G_F (determined from experiments)

Energy spectrum dependence on neutrino mass m

(from Fermi’s original article, published in German
on Zeitschrift für Physik, following rejection of the
English version by Nature)

**Measurable distortions for $m > 0$ near the end-point
(E_0 : max. allowed electron energy)**





Наблюдение на неутриното

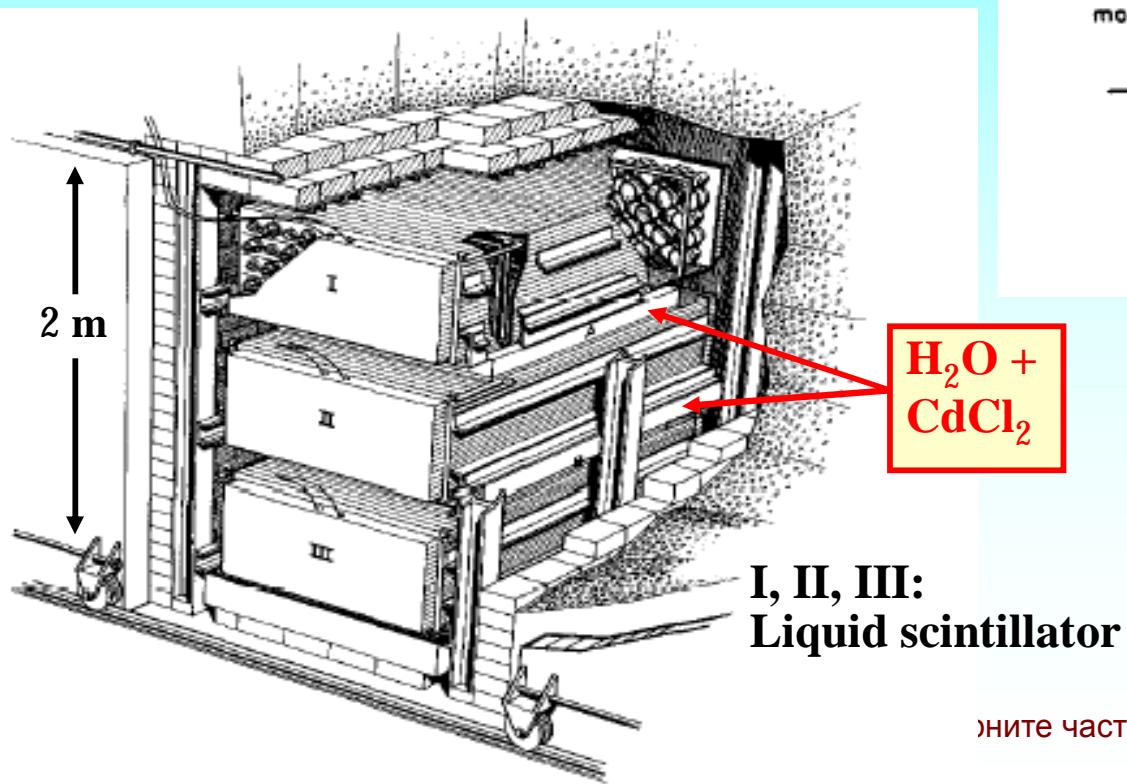


$E_g = 0.5 \text{ MeV}$

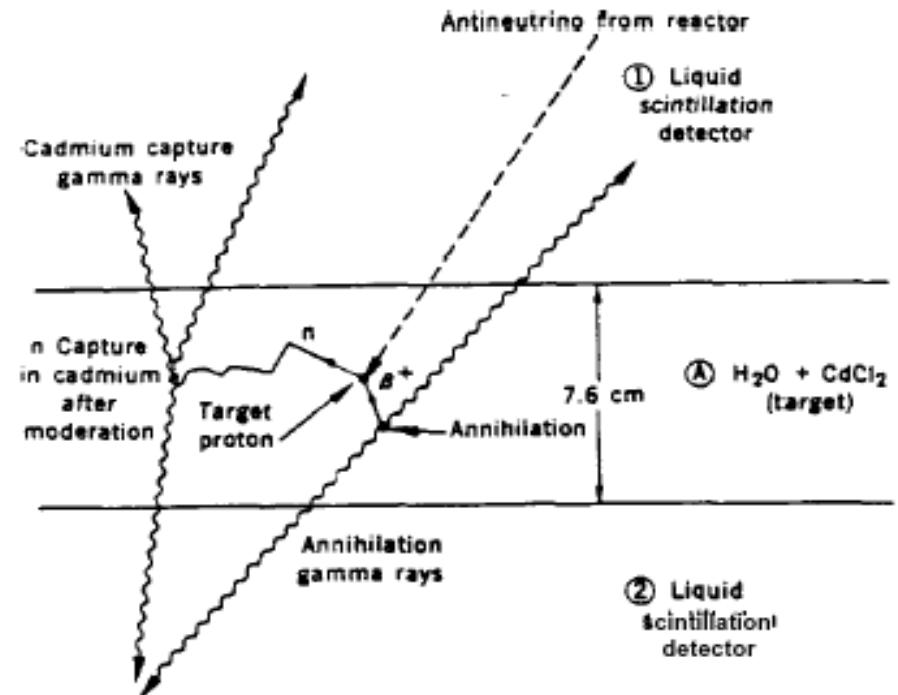


§ detect 0.5 MeV g-rays from $e^+e^- \xrightarrow{\gamma} g g$
($t = 0$)

§ neutron “thermalization” followed
by capture in Cd nuclei β emission
of delayed g-rays (average delay $\sim 30 \text{ ms}$)



юните частици



(Reines, Cowan 1953)

Event rate at the Savannah River nuclear power plant:

$3.0 \pm 0.2 \text{ events / hour}$

(after subtracting event rate measured with reactor OFF)

in agreement with expectations

Sofia, April 2009



1947: Открытие на π -мезон



Observation of the $p^+ \rightarrow m^+ \rightarrow e^+$ decay chain in nuclear emulsion exposed to cosmic rays at high altitudes

In all events the muon has a fixed kinetic energy (4.1 MeV, corresponding to a range of ~ 600 mm in nuclear emulsion) P two-body decay

$$m_p = 139.57 \text{ MeV}/c^2 ; \text{ spin} = 0$$

Dominant decay mode: $p^+ \rightarrow m^+ + n$ (and $p^- \rightarrow m^- + n\bar{\nu}$)

$$\text{Mean life at rest: } t_p = 2.6 \times 10^{-8} \text{ s} = 26 \text{ ns}$$

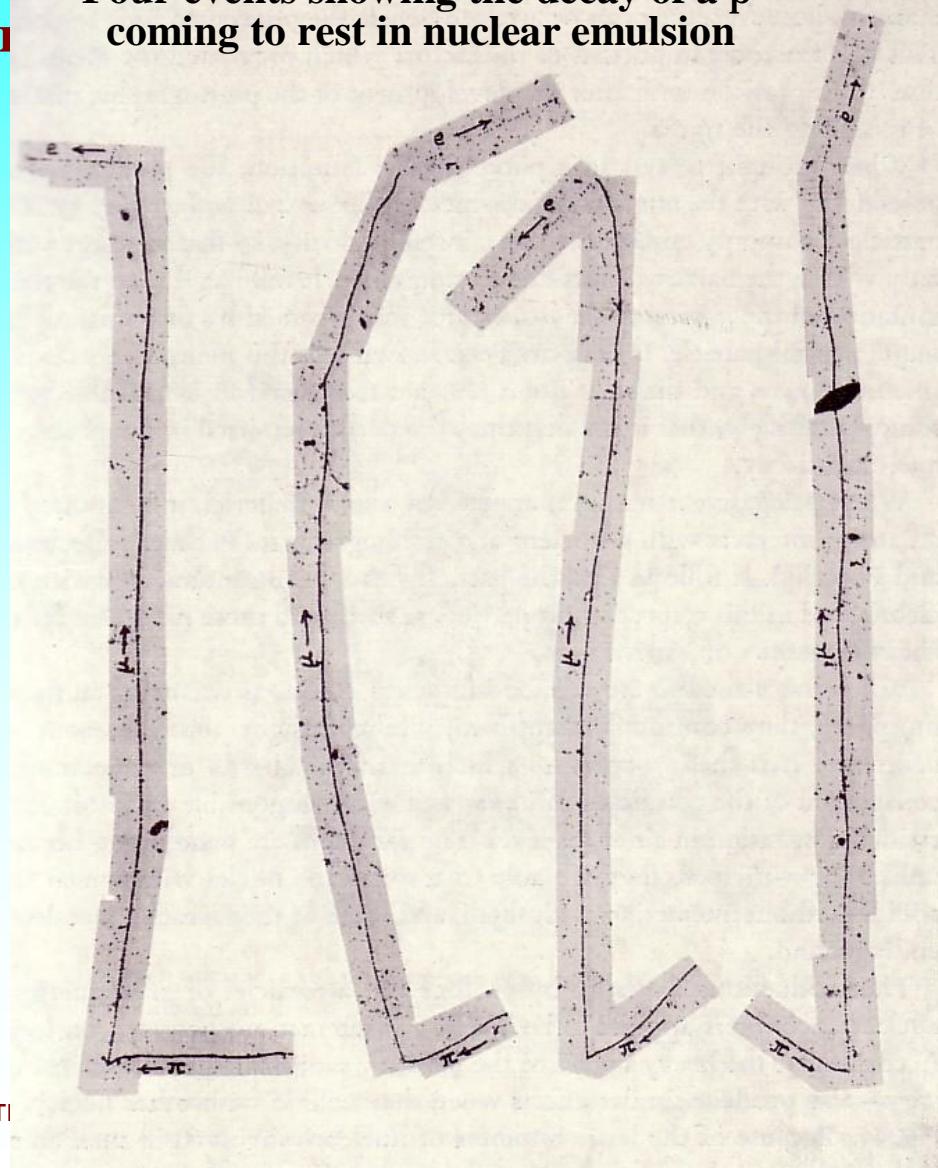
A neutral p – meson (p^0) also exists:

$$m(p^0) = 134.98 \text{ MeV}/c^2$$

Decay: $p^0 \rightarrow g + g$, mean life = $8.4 \times 10^{-17} \text{ s}$

p – mesons are the most copiously produced particles in proton – proton and proton – nucleus collisions at high energies

Four events showing the decay of a p^+ coming to rest in nuclear emulsion





Запазващи се квантови числа



Possible proton decay modes (allowed by all known conservation laws: energy – momentum, electric charge, angular momentum):

$$p \xrightarrow{\text{R}} p^0 + e^+$$

$$p \xrightarrow{\text{R}} p^0 + m^+$$

$$p \xrightarrow{\text{R}} p^+ + n$$

.....

Why is the free proton stable?

No proton decay ever observed – the proton is STABLE

Limit on the proton mean life: $t_p > 1.6 \times 10^{32}$ years

Invent a new quantum number : “Baryonic Number” B

B = 1 for proton, neutron

B = -1 for antiproton, antineutron

B = 0 for e^\pm , m^\pm , neutrinos, mesons, photons

Require conservation of baryonic number in all particle processes:

$$\sum_i B_i = \sum_f B_f$$

(i : initial state particle ; f : final state particle)



Странност



Late 1940's: discovery of a variety of heavier mesons (K – mesons) and baryons (“hyperons”) – studied in detail in the 1950's at the new high-energy proton synchrotrons (the 3 GeV “cosmotron” at the Brookhaven National Lab and the 6 GeV Bevatron at Berkeley)

Mass values

Mesons (spin = 0): $m(K^\pm) = 493.68 \text{ MeV}/c^2$; $m(K^\circ) = 497.67 \text{ MeV}/c^2$

Hyperons (spin = ½): $m(L) = 1115.7 \text{ MeV}/c^2$; $m(S^\pm) = 1189.4 \text{ MeV}/c^2$

$m(X^\circ) = 1314.8 \text{ MeV}/c^2$; $m(X^-) = 1321.3 \text{ MeV}/c^2$

Properties

§ Abundant production in proton – nucleus , p – nucleus collisions

§ Production cross-section typical of strong interactions ($S > 10^{-27} \text{ cm}^2$)

§ Production in pairs (example: $p^- + p \xrightarrow{\text{R}} K^\circ + L$; $K^- + p \xrightarrow{\text{R}} X^- + K^+$)

§ Decaying to lighter particles with mean life values $10^{-8} - 10^{-10} \text{ s}$ (as expected for a weak decay)

Examples of decay modes

$K^\pm \xrightarrow{\text{R}} p^\pm p^\circ$; $K^\pm \xrightarrow{\text{R}} p^\pm p^+ p^-$; $K^\pm \xrightarrow{\text{R}} p^\pm p^\circ p^\circ$; $K^\circ \xrightarrow{\text{R}} p^+ p^-$; $K^\circ \xrightarrow{\text{R}} p^\circ p^\circ$; ...

$L \xrightarrow{\text{R}} p p^-$; $L \xrightarrow{\text{R}} n p^\circ$; $S^+ \xrightarrow{\text{R}} p p^\circ$; $S^+ \xrightarrow{\text{R}} n p^+$; $S^+ \xrightarrow{\text{R}} n p^-$; ...

$X^- \xrightarrow{\text{R}} L p^-$; $X^\circ \xrightarrow{\text{R}} L p^\circ$



Странност



(Gell-Mann, Nakano, Nishijima, 1953)

§ conserved in strong interaction processes:

$$\sum_i S_i = \sum_f S_f$$

§ not conserved in weak decays:

$$\left| S_i - \sum_f S_f \right| = 1$$

$S = +1$: K^+, K° ; $S = -1$: L, S^\pm, S° ; $S = -2$: X°, X^- ; $S = 0$: all other particles
(and opposite strangeness $-S$ for the corresponding antiparticles)

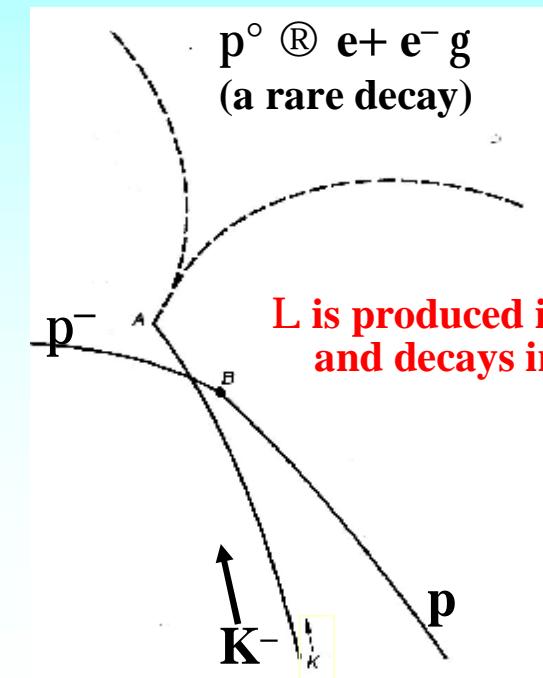
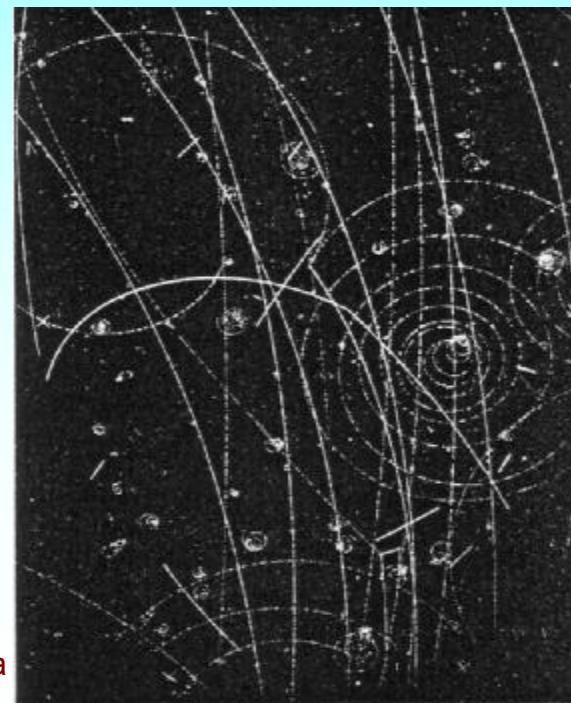
Example of a K^- stopping
in liquid hydrogen:



(strangeness conserving)
followed by the decay



(strangeness violation)



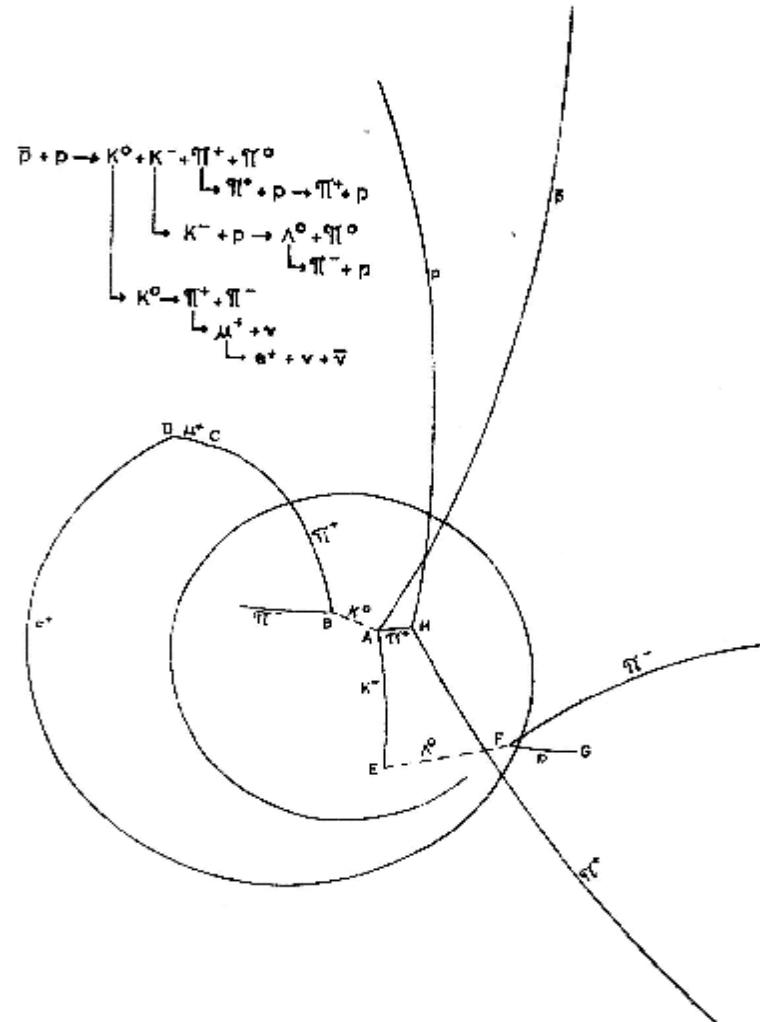


Откритие на антипротона



L. Litov

Физика на елементарните частици



Sofia, April 2009



A puzzle of the late 1950's: the absence of $m^+ \rightarrow e^+ \gamma$ decays

Experimental limit: < 1 in 10^6 $m^+ \rightarrow e^+ n \bar{n}$ decays

A possible solution: existence of a new, conserved “muonic” quantum number distinguishing muons from electrons

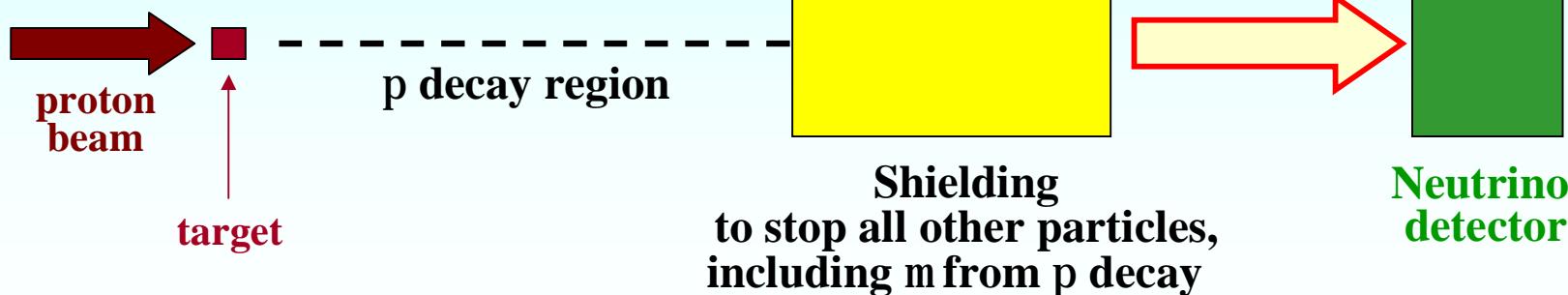
To allow $m^+ \rightarrow e^+ n \bar{n}$ decays, \bar{n} must have “muonic” quantum number but not n . In m^+ decay the \bar{n} is not the antiparticle of n

Two distinct neutrinos (n_e , n_m) in the decay $m^+ \rightarrow e^+ n_e \bar{n}_m$

Consequence for p – meson decays: $p^+ \rightarrow m^+ n_m$; $p^- \rightarrow m^- \bar{n}_m$ to conserve the “muonic” quantum number

High energy proton accelerators: intense sources of p^\pm – mesons $\rightarrow n_m, \bar{n}_m$

Experimental method



If $n_m + n_e \rightarrow n_m$ interactions produce m^- and not e^- (example: $n_m + n \rightarrow m^- + p$)



Кварков модел



Late 1950's – early 1960's: discovery of many strongly interacting particles at the high energy proton accelerators (Berkeley Bevatron, BNL AGS, CERN PS), all with very short mean life times ($10^{-20} - 10^{-23}$ s, typical of strong decays)
P catalog of > 100 strongly interacting particles (collectively named "hadrons")

ARE HADRONS ELEMENTARY PARTICLES?

1964 (Gell-Mann, Zweig): Hadron classification into "families";
observation that all hadrons could be built from three spin $\frac{1}{2}$
"building blocks" (named "quarks" by Gell-Mann):

	<i>u</i>	<i>d</i>	<i>s</i>
Electric charge (units $ e $)	+2/3	- 1/3	- 1/3
Baryonic number	1/3	1/3	1/3
Strangeness	0	0	- 1

and three antiquarks (\bar{u} , \bar{d} , \bar{s}) with opposite electric charge
and opposite baryonic number and strangeness

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Кварков модел



Mesons: quark – antiquark pairs

Examples of non-strange mesons:

$$p^+ \equiv u\bar{d} \quad ; \quad p^- \equiv \bar{u}d \quad ; \quad p^0 \equiv (d\bar{d} - u\bar{u})/\sqrt{2}$$

Examples of strange mesons:

$$K^- \equiv s\bar{u} \quad ; \quad \bar{K}^0 \equiv s\bar{d} \quad ; \quad K^+ \equiv \bar{s}u \quad ; \quad K^0 \equiv \bar{s}d$$

Baryons: three quarks bound together

Antibaryons: three antiquarks bound together

Examples of non-strange baryons:

$$\text{proton} \equiv uud \quad ; \quad \text{neutron} \equiv udd$$

Examples of strangeness -1 baryons:

$$\Sigma^+ \equiv suu \quad ; \quad \Sigma^0 \equiv sud \quad ; \quad \Sigma^- \equiv sdd$$

Examples of strangeness -2 baryons:

$$\Xi^0 \equiv ssu \quad ; \quad \Xi^- \equiv ssd$$

Физика на елементарните частици



A success of the static quark model



The “decuplet” of spin $\frac{3}{2}$ baryons

Strangeness					Mass (MeV/c ²)
0	N^{*++} <i>uuu</i>	N^{*+} <i>uud</i>	$N^{*\circ}$ <i>udd</i>	N^{*-} <i>ddd</i>	1232
-1	S^{*+} <i>suu</i>	$S^{*\circ}$ <i>sud</i>	S^{*-} <i>sdd</i>		1384
-2		$X^{*\circ}$ <i>ssu</i>	X^{*-} <i>ssd</i>		1533
-3			W^- <i>sss</i>		1672 (predicted)

W^- : the bound state of three s – quarks with the lowest mass
with total angular momentum = $3/2$ P



Pauli's exclusion principle requires that the three quarks
cannot be identical

A success of the static quark model



The first W^- event (observed in the 2 m liquid hydrogen bubble chamber at BNL using a 5 GeV/c K^- beam from the 30 GeV AGS)

Chain of events in the picture:

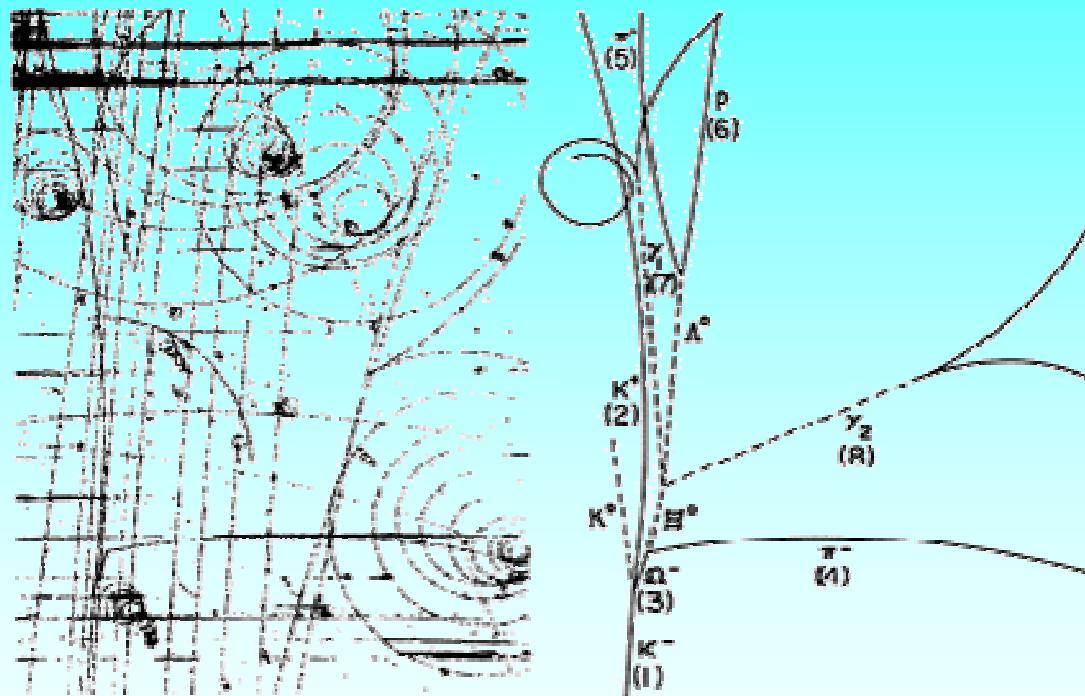
$K^- + p \xrightarrow{\text{R}} W^- + K^+ + K^0$
 (strangeness conserving)

$W^- \xrightarrow{\text{R}} X^0 + p^-$
 (DS = 1 weak decay)

$X^0 \xrightarrow{\text{R}} p^0 + L$
 (DS = 1 weak decay)

$L \xrightarrow{\text{R}} p^- + p$
 (DS = 1 weak decay)

$p^0 \xrightarrow{\text{R}} g + g$ (electromagnetic decay)
 with both g – rays converting to an e^+e^- in liquid hydrogen
 (very lucky event, because the mean free path for $g \xrightarrow{\text{R}} e^+e^-$ in liquid hydrogen is ~ 10 m)



W^- mass measured from this event = 1686 ± 12 MeV/ c^2



Дълбоко нееластично разсейване

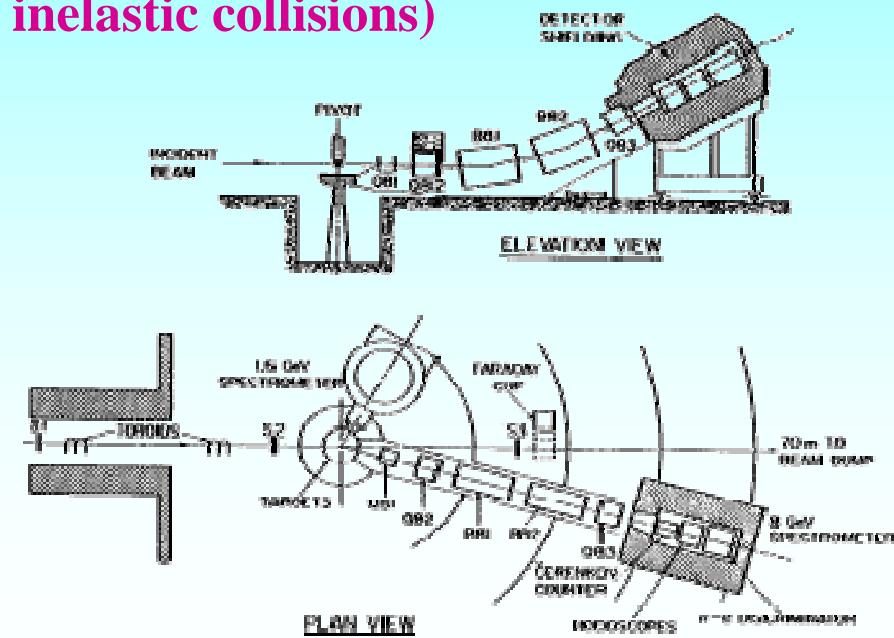
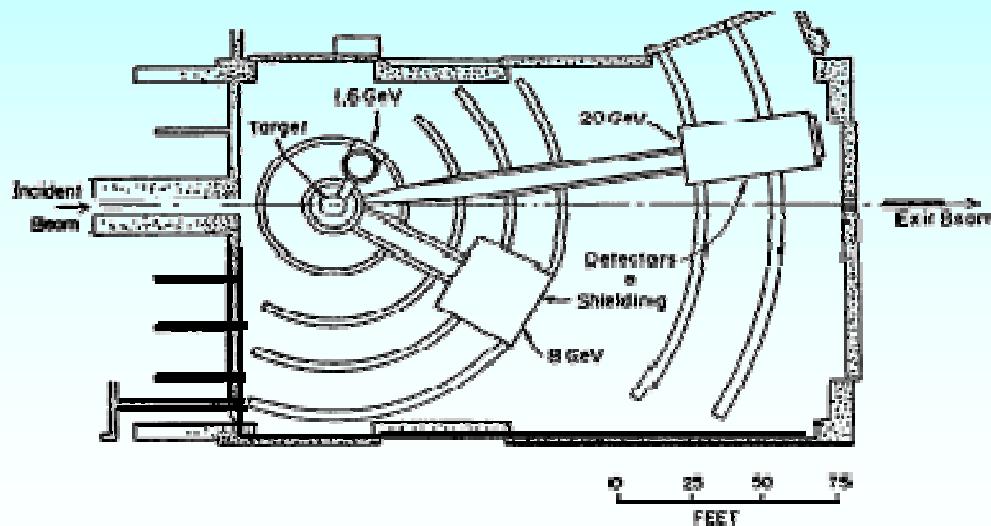


Electron – proton scattering using a 20 GeV electron beam from the Stanford two – mile Linear Accelerator (1968 – 69).

**The modern version of Rutherford's original experiment:
resolving power » wavelength associated with 20 GeV electron » 10^{-15} cm**

Three magnetic spectrometers to detect the scattered electron:

- § 20 GeV spectrometer (to study elastic scattering $e^- + p \xrightarrow{\text{R}} e^- + p$)
- § 8 GeV spectrometer (to study inelastic scattering $e^- + p \xrightarrow{\text{R}} e^- + \text{hadrons}$)
- § 1.6 GeV spectrometer (to study extremely inelastic collisions)





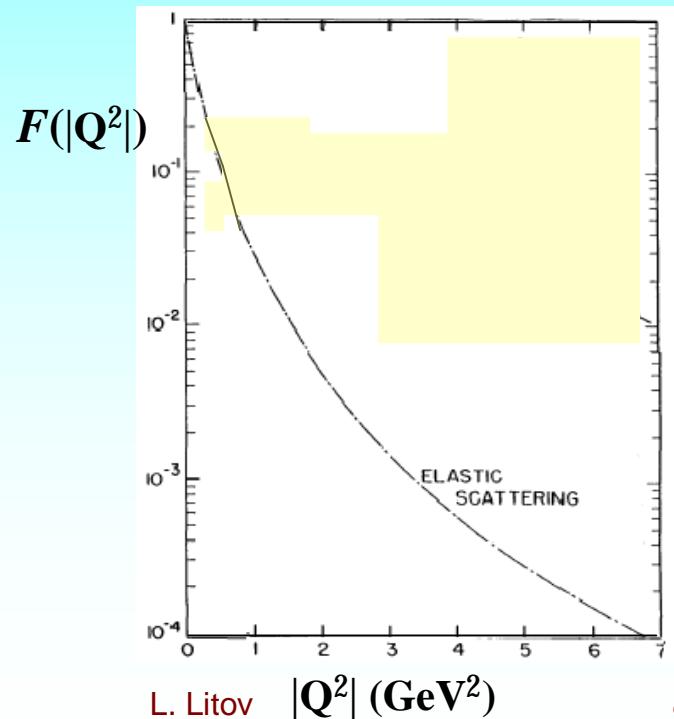
Дълбоко нееластично разсейване



Electron elastic scattering from a point-like charge $|e|$ at high energies:
differential cross-section in the collision centre-of-mass (Mott's formula)

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2 (\hbar c)^2}{8 E^2} \frac{\cos^2(\theta/2)}{\sin^4(\theta/2)} \equiv \sigma_M \quad \left(\alpha = \frac{e^2}{\hbar c} \approx \frac{1}{137} \right)$$

Scattering from an extended charge distribution: multiply σ_M by a “form factor”:



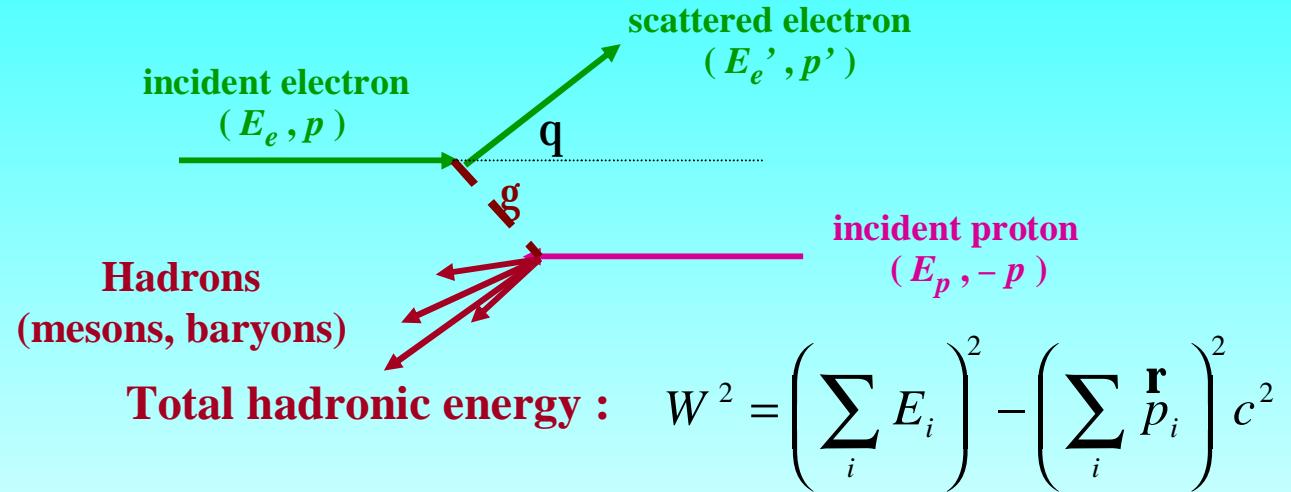
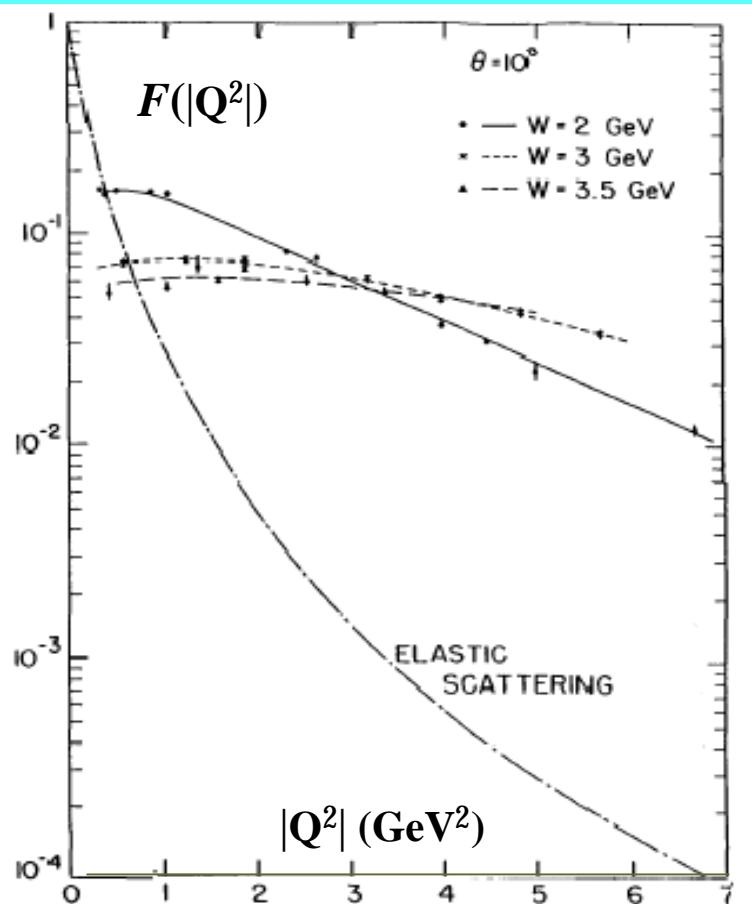
$$\frac{d\sigma}{d\Omega} = F(|Q^2|) \sigma_M$$

$|Q| = \hbar / D$: mass of the exchanged virtual photon
D: linear size of target region contributing to scattering
Increasing $|Q| \rightarrow$ decreasing target electric charge

$F(|Q^2|) = 1$ for a point-like particle
→ the proton is not a point-like particle



Дълбоко нееластично разсейване



For deeply inelastic collisions,
the cross-section depends only weakly
on $|Q^2|$, suggesting a collision
with a POINT-LIKE object



Партонаен модел



Deep inelastic electron – proton collisions are elastic collisions with point-like, electrically charged, spin $\frac{1}{2}$ constituents of the proton carrying a fraction x of the incident proton momentum

Each constituent type is described by its electric charge e_i (units of $|e|$) and by its x distribution (dN_i/dx) (“structure function”)

If these constituents are the u and d quarks, then deep inelastic e – p collisions provide information on a particular combination of structure functions:

$$\left(\frac{dN}{dx} \right)_{e-p} = e_u^2 \frac{dN_u}{dx} + e_d^2 \frac{dN_d}{dx}$$

Comparison with $n_m - p$ and $\bar{n}_m - p$ deep inelastic collisions at high energies under the assumption that these collisions are also elastic scatterings on quarks

$n_m + p \xrightarrow{\text{R}} m^- + \text{hadrons} : n_m + d \xrightarrow{\text{R}} m^- + u$ (depends on dN_d/dx)

$\bar{n}_m + p \xrightarrow{\text{R}} m^+ + \text{hadrons} : n_m \bar{n} + u \xrightarrow{\text{R}} m^+ + d$ (depends on dN_u/dx)

(Neutrino interactions do not depend on electric charge)

All experimental results on deep inelastic $e - p$, $n_m - p$, $n_m - \bar{p}$ collisions are consistent with $e_u^2 = 4/9$ and $e_d^2 = 1/9$

→ the proton constituents are the quarks



Цвят



Problem with

J=3/2 baryon $\Delta^{++} \Rightarrow u \uparrow u \uparrow u \uparrow$

It has symmetric wave function, but it is a fermion

Contradiction with Pauli exclusion principle –
the wave function should be antisymmetric

The way out - new quantum number – colour

$N_c = 3$ q^α , $\alpha = 1, 2, 3$

Then $\Delta^{++} = \frac{1}{\sqrt{6}} e^{abg} u_a \uparrow u_b \uparrow u_g \uparrow$

In general case

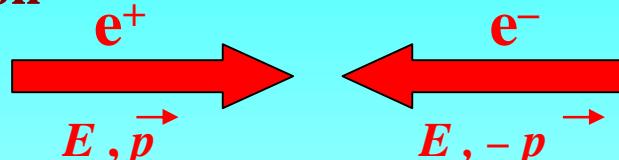
$$B = \frac{1}{\sqrt{6}} e^{abg} | q_a q_b q_g \rangle \quad M = \frac{1}{\sqrt{3}} d^{ab} | q_a \bar{q}_b \rangle$$



Експериментална проверка на хипотезата за цвят



Two beams circulating in opposite directions in the same magnetic ring and colliding head-on



A two-step process: $e^+ + e^- \rightarrow$ virtual photon $\rightarrow f + \bar{f}$

f : any electrically charged elementary spin $\frac{1}{2}$ particle (m , quark)
(excluding e^+e^- elastic scattering)

Virtual photon energy – momentum : $E_g = 2E, p_g = 0$ $\Rightarrow Q^2 = E_g^2 - p_g^2 c^2 = 4E^2$

Cross - section for $e^+e^- \rightarrow f\bar{f}$: $S = \frac{2pa^2 h^2 c^2}{3Q^2} e_f^2 b(3-b)$

$$a = e^2/(\hbar c) \gg 1/137$$

e_f : electric charge of particle f (units $|e|$)

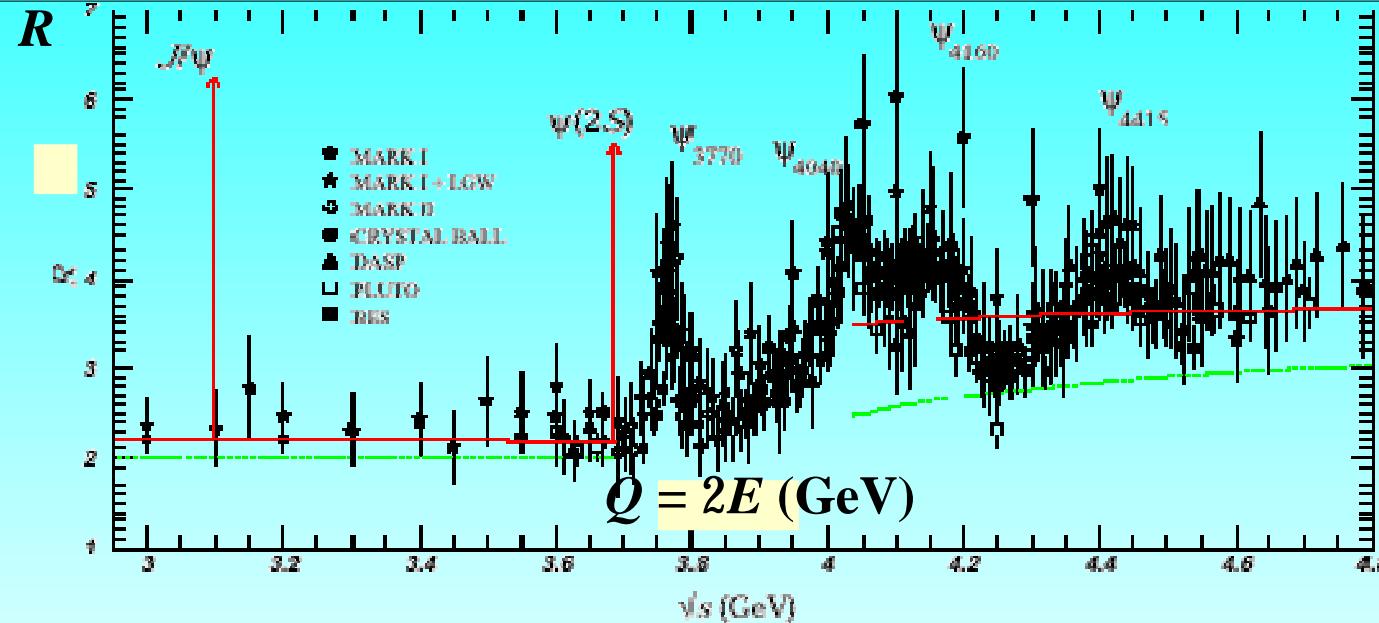
$b = v/c$ of outgoing particle f

(formula precisely verified for $e^+e^- \rightarrow m^+m^-$)

Assumption: $e^+e^- \rightarrow$ quark (q) + antiquark (\bar{q}) \rightarrow hadrons
 \Rightarrow at energies $E \gg m_q c^2$ (for $q = u, d, s$) $b \gg 1$:

$$R \equiv \frac{S(e^+e^- \rightarrow \text{hadrons})}{S(e^+e^- \rightarrow m^+m^-)} = e_u^2 + e_d^2 + e_s^2 = \frac{4}{9} + \frac{1}{9} + \frac{1}{9} = \frac{2}{3}$$

Stanford e^+e^- collider SPEAR (1974 –75):



§ For $Q < 3.6$ GeV $R \gg 2$. If each quark exists in three different states, $R \gg 2$ is consistent with $3 \times (2/3)$. This would solve the W^- problem.

§ Between 3 and 4.5 GeV, the peaks and structures are due to the production of quark-antiquark bound states and resonances of a fourth quark (“charm”, c) of electric charge $+2/3$

§ Above 4.6 GeV $R \gg 4.3$. Expect $R \gg 2$ (from u, d, s) $+ 3 \times (4/9) = 3.3$ from the addition of the c quark alone. So the data suggest pair production of an additional elementary spin $1/2$ particle with electric charge = 1 (later identified as the t – lepton (no strong interaction) with mass $\gg 1777$ MeV/ c^2).



Фундаментални частици



Leptons

	Electric Charge
Tau	-1
Muon	-1
Electron	-1
Tau Neutrino	0
Muon Neutrino	0
Electron Neutrino	0

Quarks

	Electric Charge
Bottom	-1/3
Strange	-1/3
Down	-1/3
Top	2/3
Charm	2/3
Up	2/3

each quark: ● R, ● B, ● G 3 colors

The particle drawings are simple artistic representations